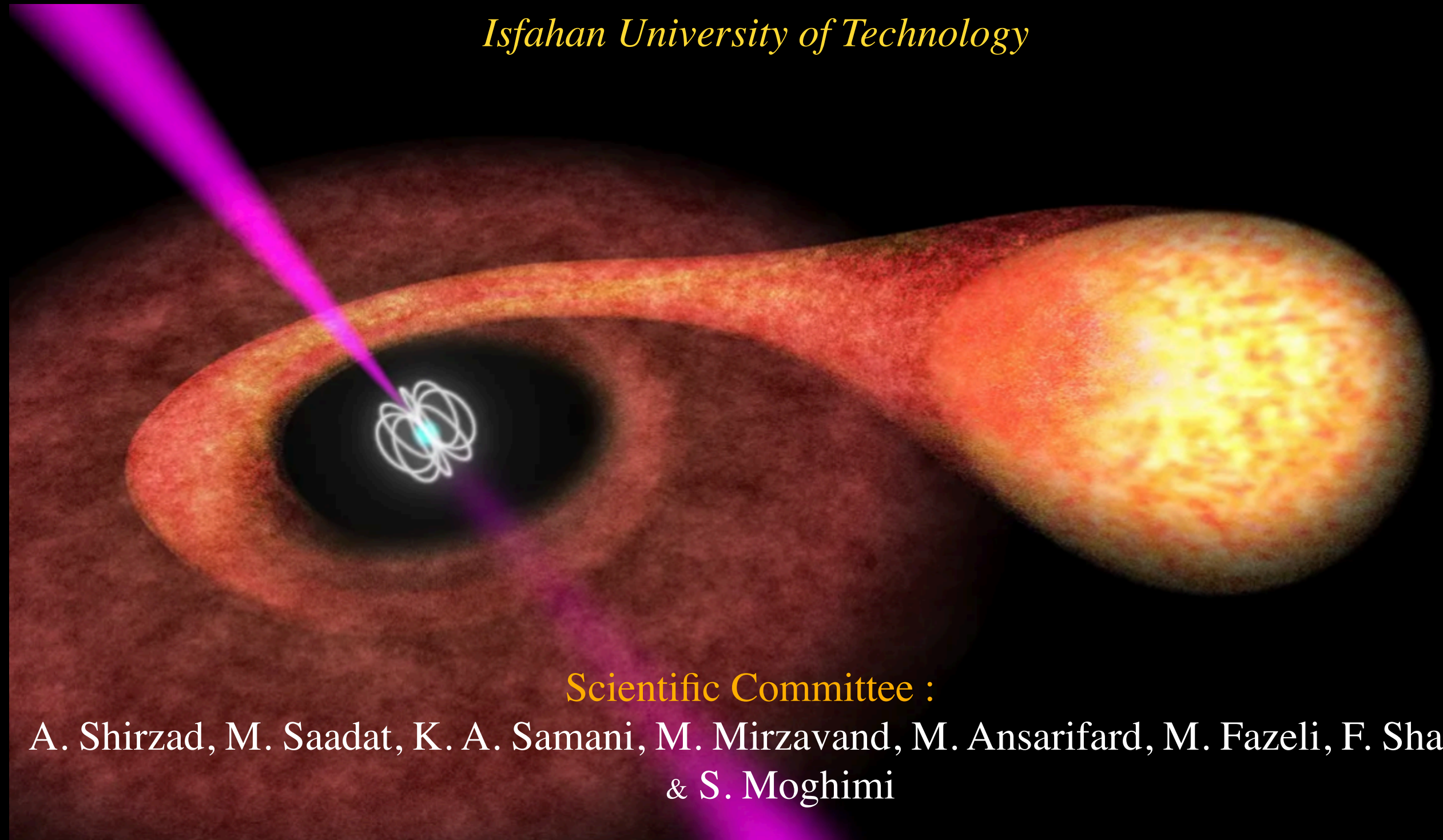


Black Widow Pulsar

The 54th International Physics Olympiad (IPhO 2024)

Soroush Shakeri

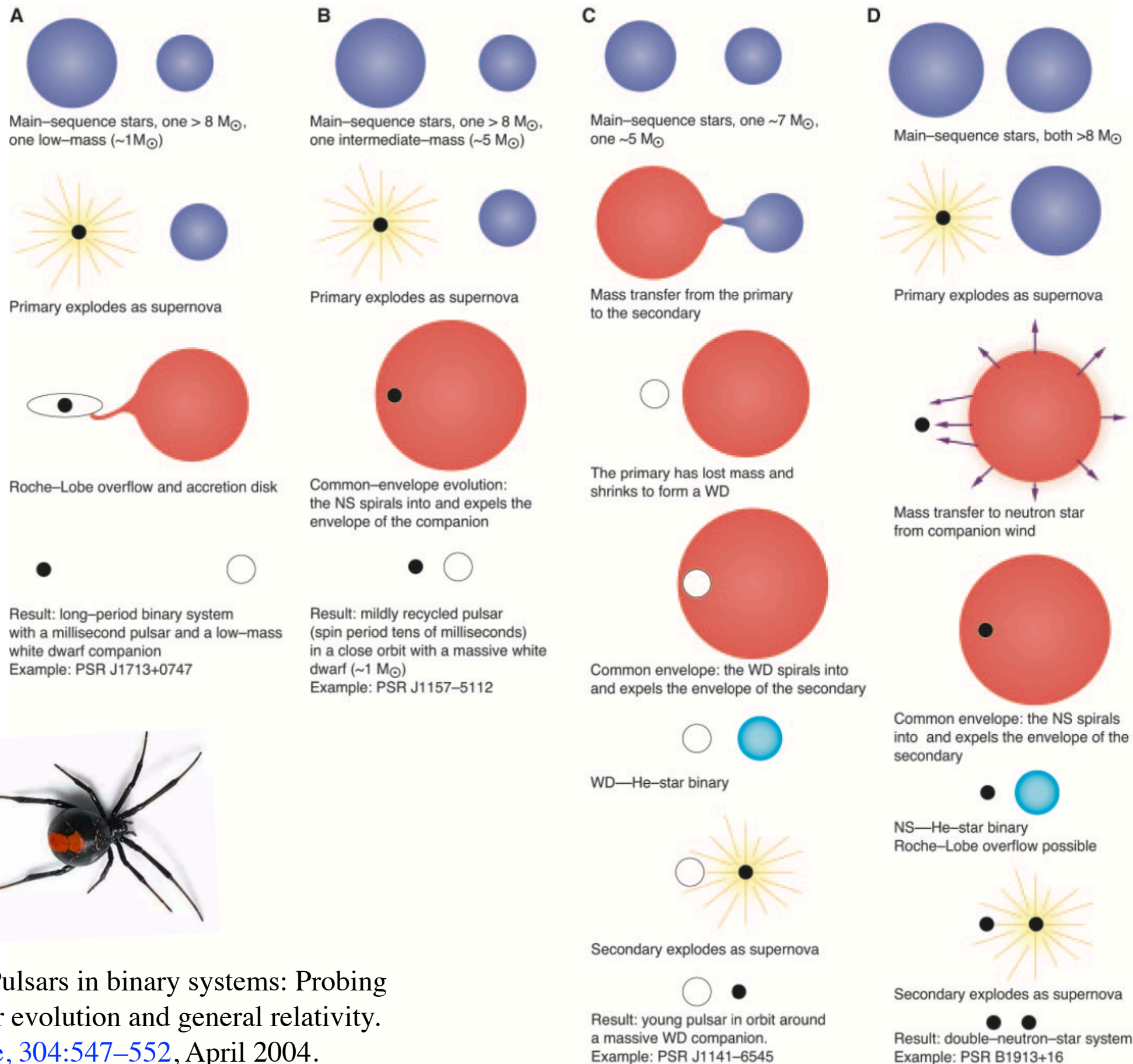
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& S. Moghimi

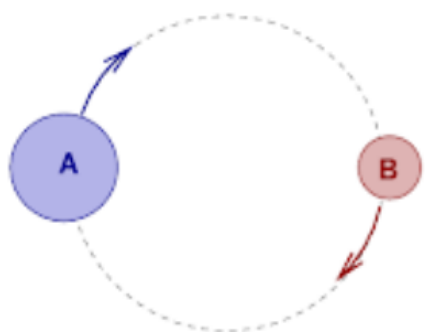
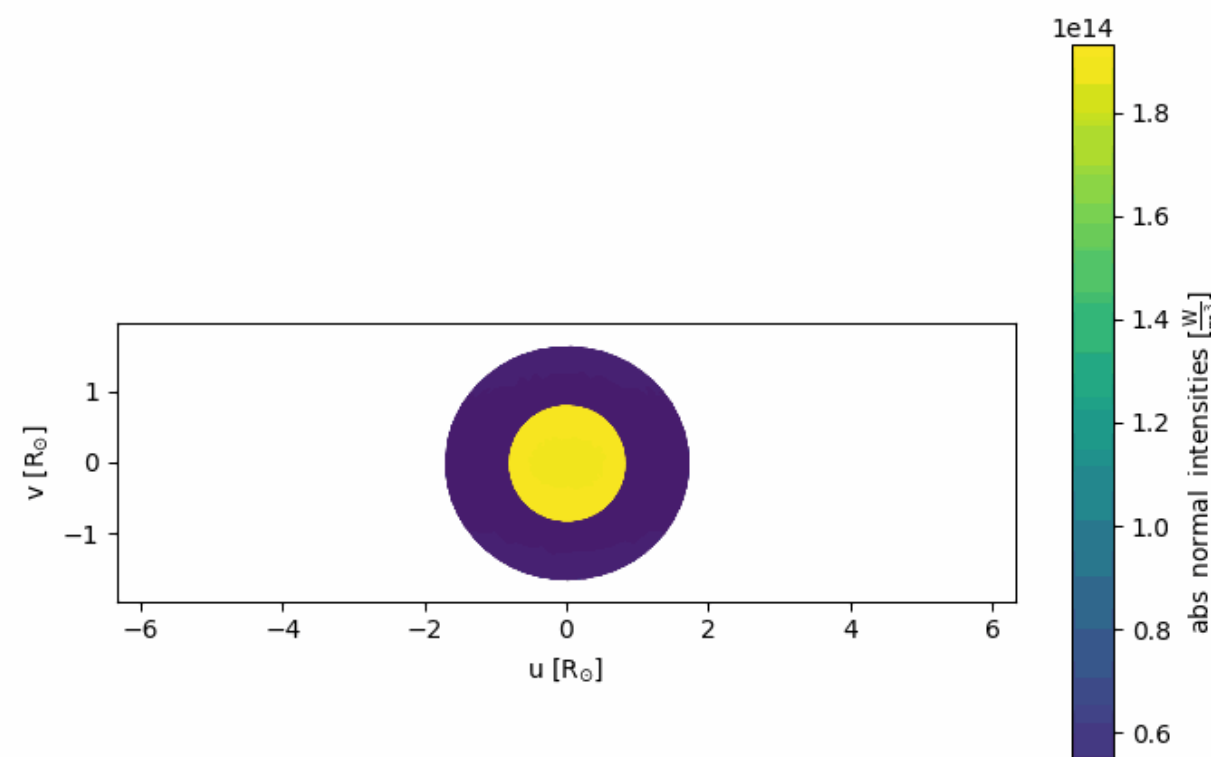
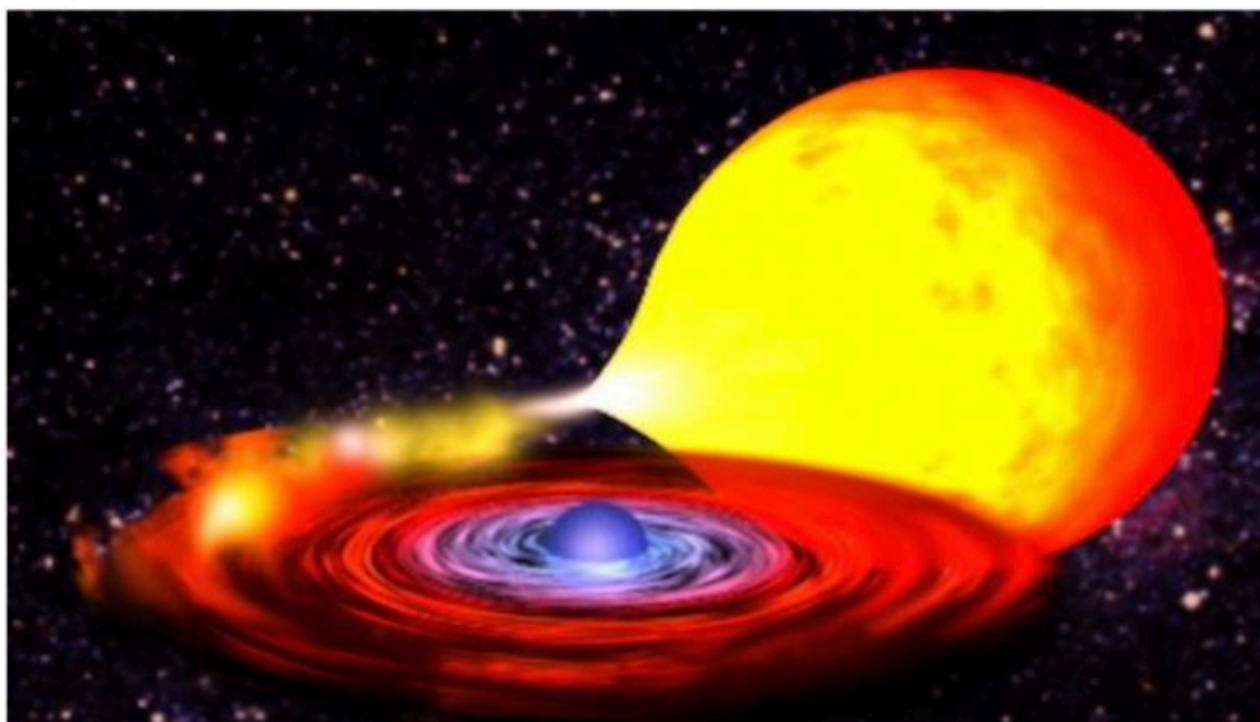
Possible Evolutionary Paths of Binary Systems



* I. H. Stairs. Pulsars in binary systems: Probing binary stellar evolution and general relativity.
[Science, 304:547–552, April 2004.](#)



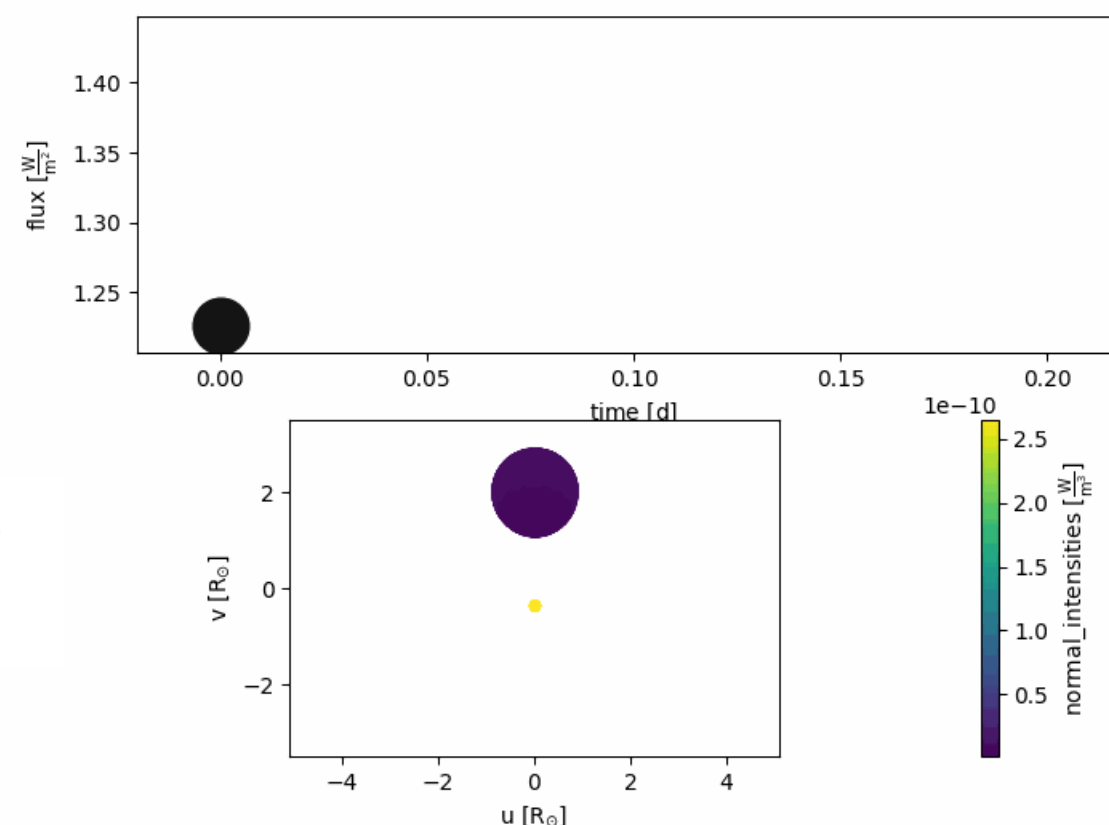
Black Widow in Eclipsing Binary System



Top view

Black Widow (Red Back) Pulsars are the heaviest Neutron star ever!

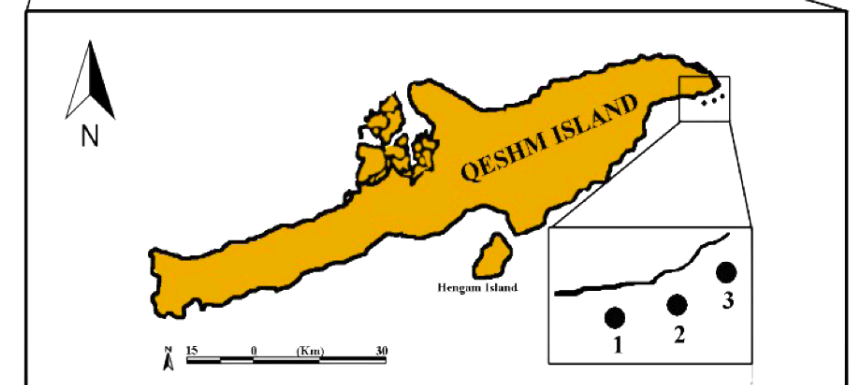
1988: Fruchter, Stinebring & Taylor (Nature 333, 237, 1988) found an eclipsing pulsar with a very low mass companion, the hypothesis of **ablation** wind quickly follows



* Fruchter, A. S.; Stinebring, D. R.; Taylor, J. H. (May 19, 1988). "A millisecond pulsar in an eclipsing binary". *Nature*. **333** (6170):

Black Widow (Red Back) Spider

A neutron star in this close binary systems has been compared to a black widow (red back) spider, (a female spider which eats its mate after mating).



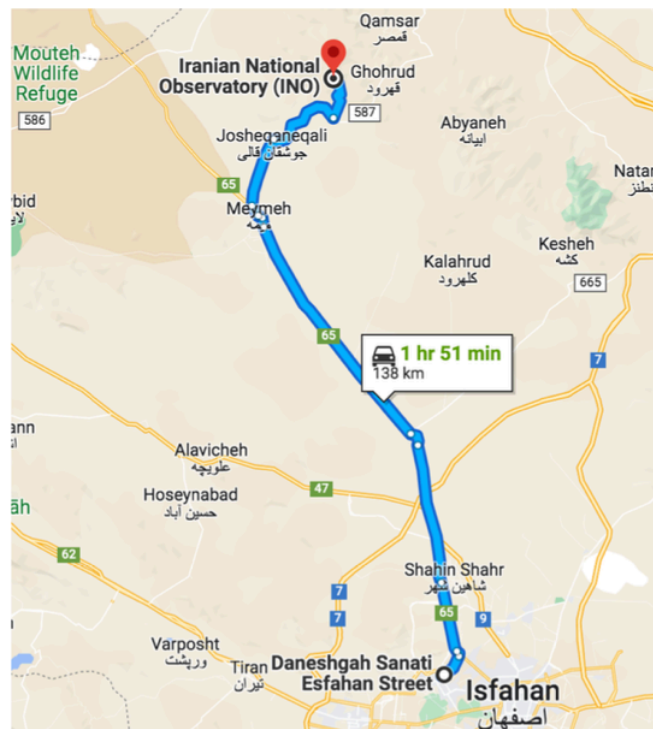
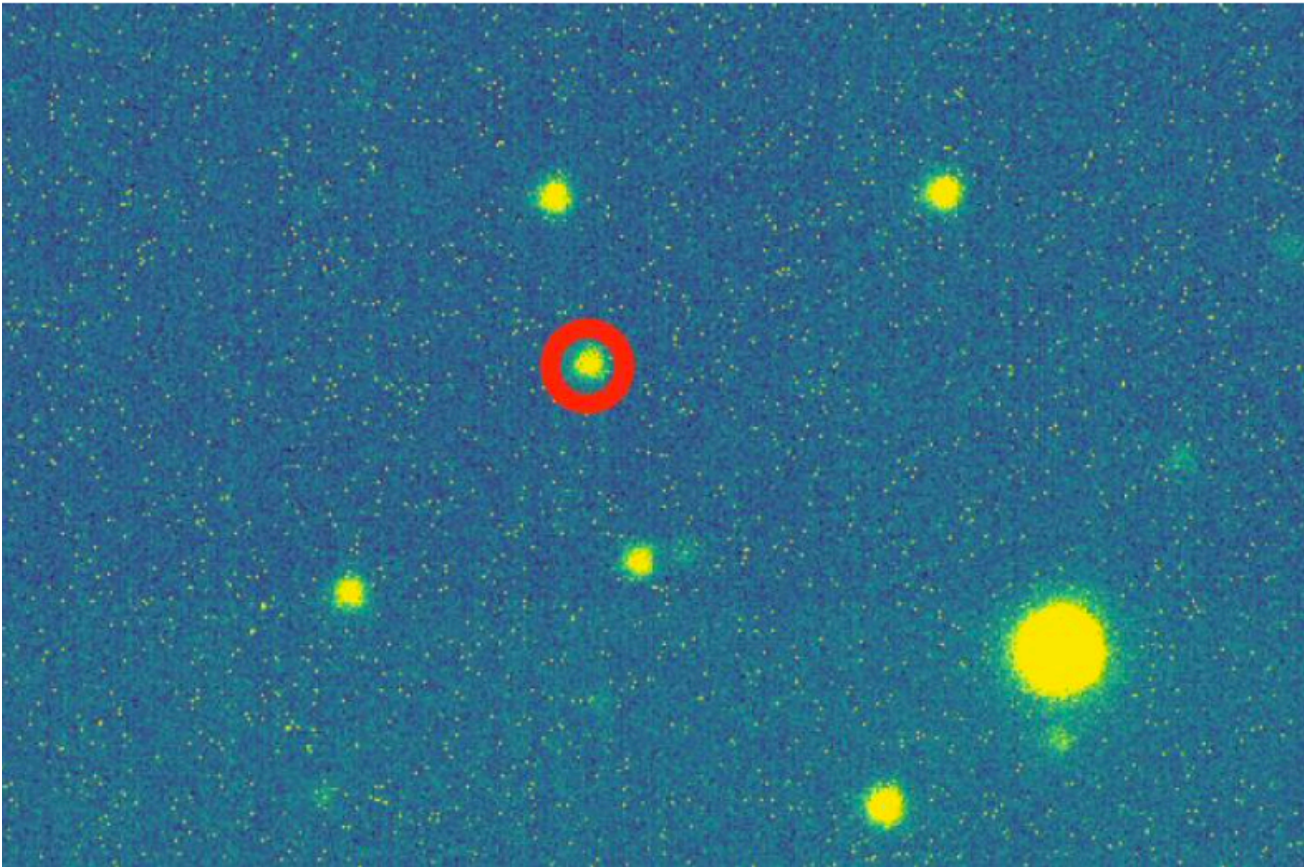
Qeshm_Dulab

July 2024 - Gheshm Island

M.eslamidulabi

Optical Observation of companion of a spider pulsar by INO

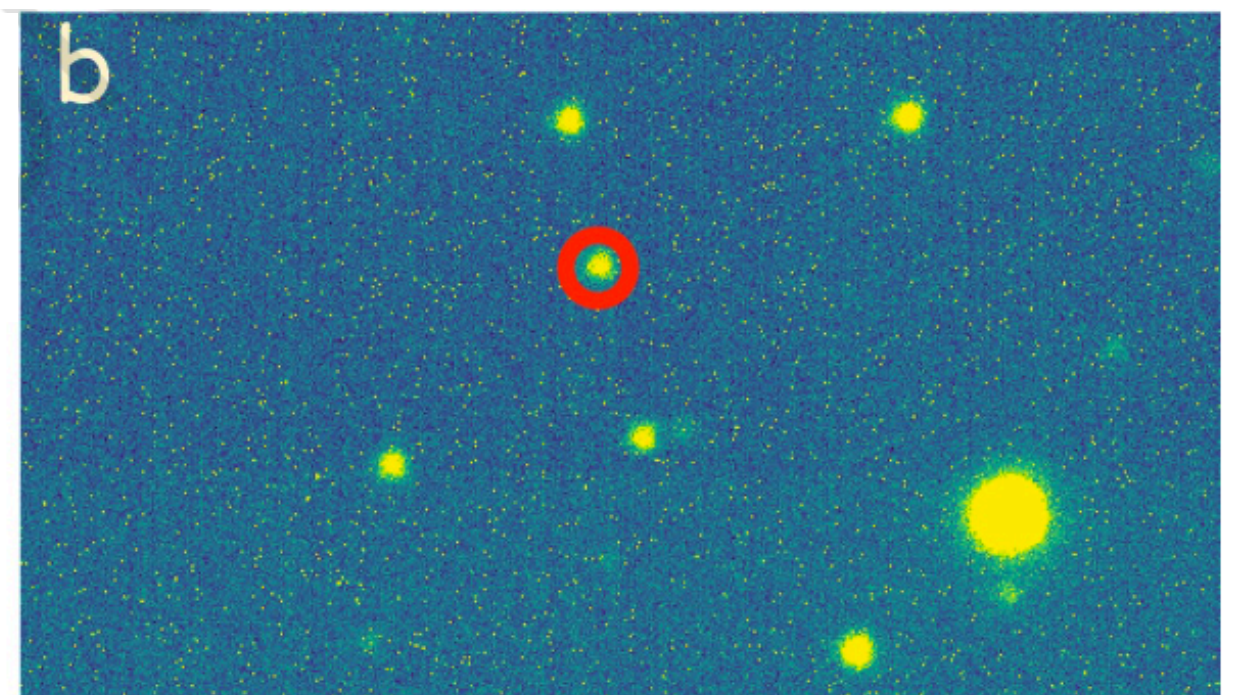
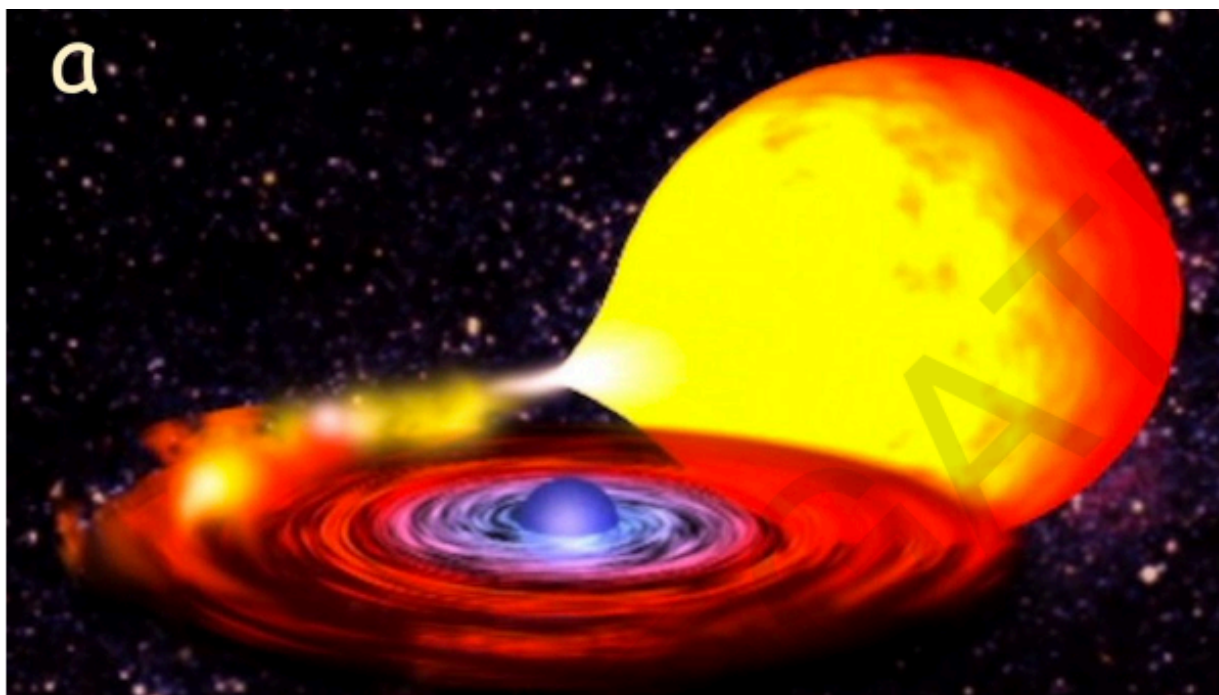
The figure below shows the picture of the companion of the neutron star PSR J2215+5135, taken by the 3.4-meter optical telescope of the Iranian National Observatory. No neutron star can be seen in this image and the observed light is due to its companion.





Black Widow Pulsars

A significant number of the observed stars are binaries. One or both of the stars may be neutron stars rotating with a high angular velocity and emitting electromagnetic waves; such stars are called pulsars. Sometimes a companion star is an expansive mass of gas that gradually falls down onto the neutron star and causes its mass to increase (Figure 1-a). In this way, a neutron star gradually swallows up a portion of the mass of its companion star. For this reason, the neutron star has been compared to a black widow (or redback spider), a female spider which eats its mate after mating. The heating of the gas falling down onto the black widow generates radiation which can be observed. The heaviest neutron stars often are black widows and they serve as natural laboratories for testing fundamental physics. Figure 1-b shows the picture of the companion of the neutron star PSR J2215+5135, taken by the 3.4-meter optical telescope of the Iranian National Observatory. No neutron star can be seen in this image and the observed light is due to its companion.

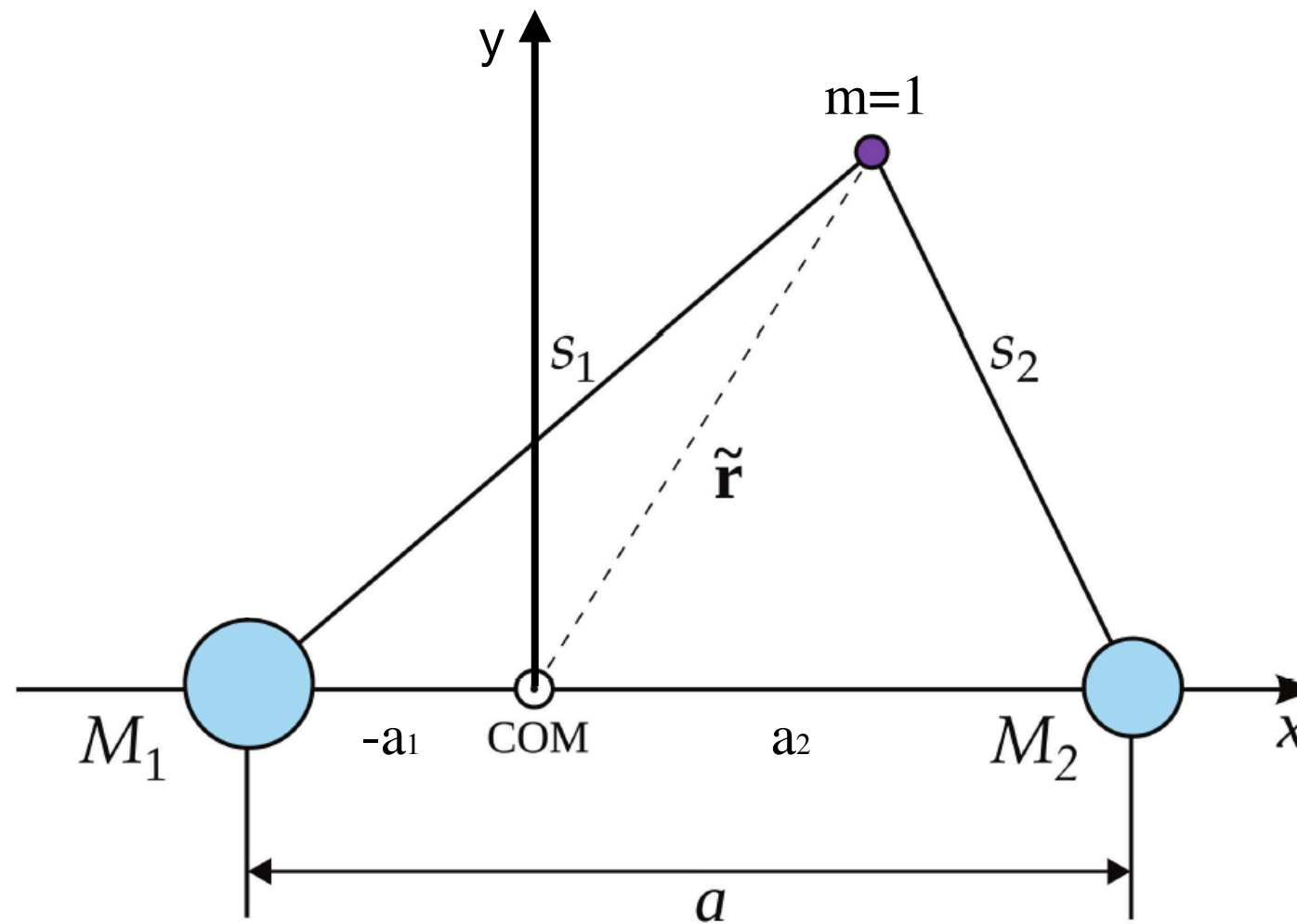


Question

Consider a simple model in which the black widow and its companion star, are represented by two point masses M_1 and M_2 moving on a circular orbit around their center of mass. To investigate the dynamics of this system, consider a rotating coordinate system in which the two bodies are stationary. Take the center of mass to be the origin of the coordinate system. Assume that the two point masses lie on the x -axis on both sides of the origin at a distance a from each other, and that M_1 lies on the negative x -axis. At an arbitrary point (x, y) in the plane of motion, the effective potential $\varphi(x, y)$ for a unit test mass is the sum of the gravitational potentials of the two point masses plus the centrifugal potential.

A-1	Write $\varphi(x, y)$ in terms of M_1 , M_2 , and a .	1.0 pt
A-2	Assuming $M_1 > M_2$, plot the function $\varphi(x, 0)$ qualitatively.	0.7 pt

Solution A.1.



$$\begin{cases} M_1 a_1 = M_2 a_2 \\ a_1 + a_2 = a \end{cases} \rightarrow a_1 = \frac{M_2}{M} a, a_2 = \frac{M_1}{M} a : M = M_1 + M_2$$

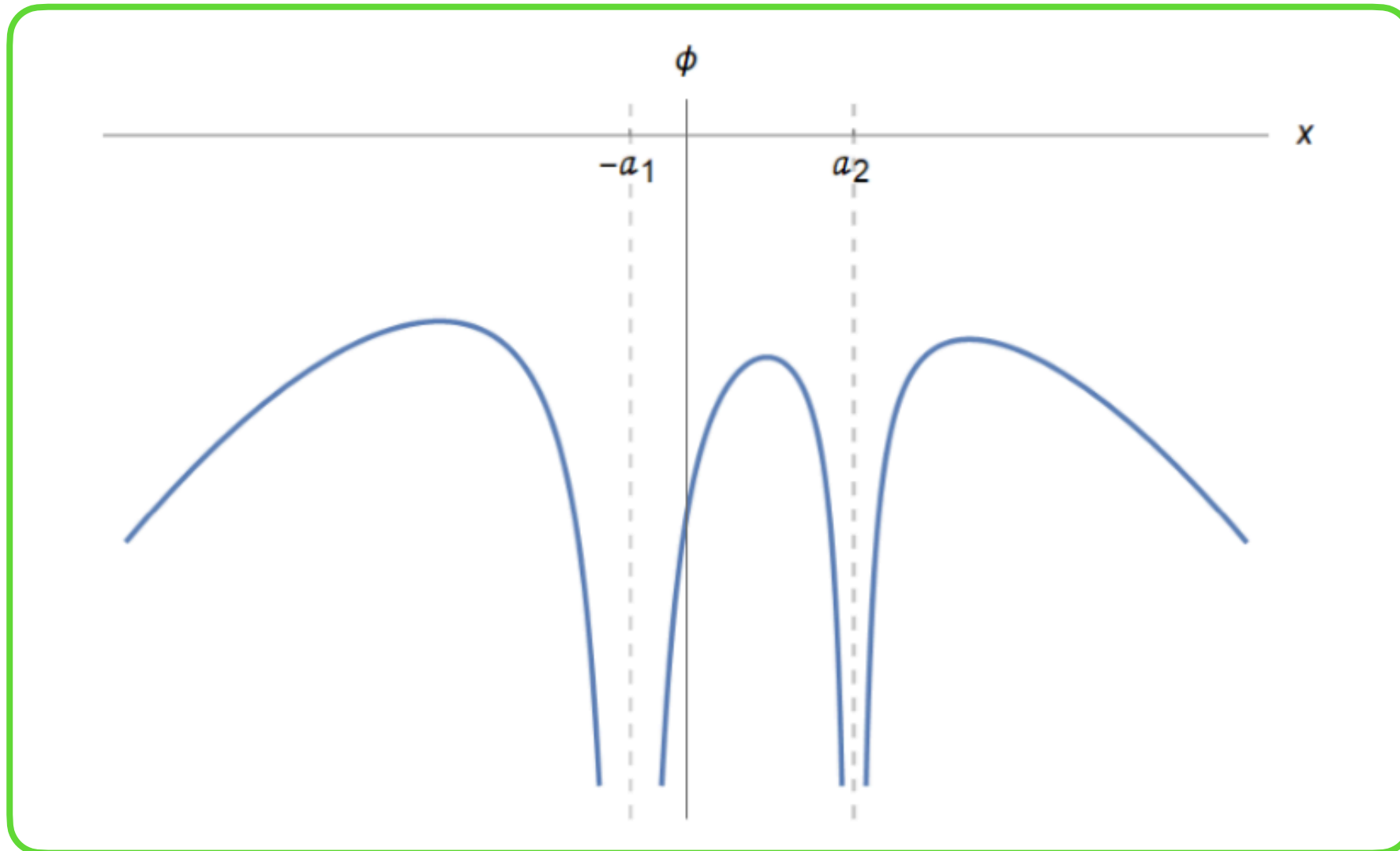
$$U = -\frac{1}{2} \omega^2 r^2 \quad \omega = \sqrt{\frac{GM}{a^3}}$$

$$\Phi(x, y) = -\frac{GM_1}{s_1} - \frac{GM_2}{s_2} - \frac{1}{2} \omega^2 \tilde{r}^2$$

$$\Phi(x, y) = -\frac{GM_1}{\sqrt{(x + a_1)^2 + y^2}} - \frac{GM_2}{\sqrt{(x - a_2)^2 + y^2}} - \frac{1}{2} \omega^2 (x^2 + y^2)$$

Solution A.2.

$$\Phi(x, 0) = -\frac{GM_1}{\left|x + \frac{M_2}{M}a\right|} - \frac{GM_2}{\left|x - \frac{M_1}{M}a\right|} - \frac{1}{2} \frac{GM}{a^3} x^2$$

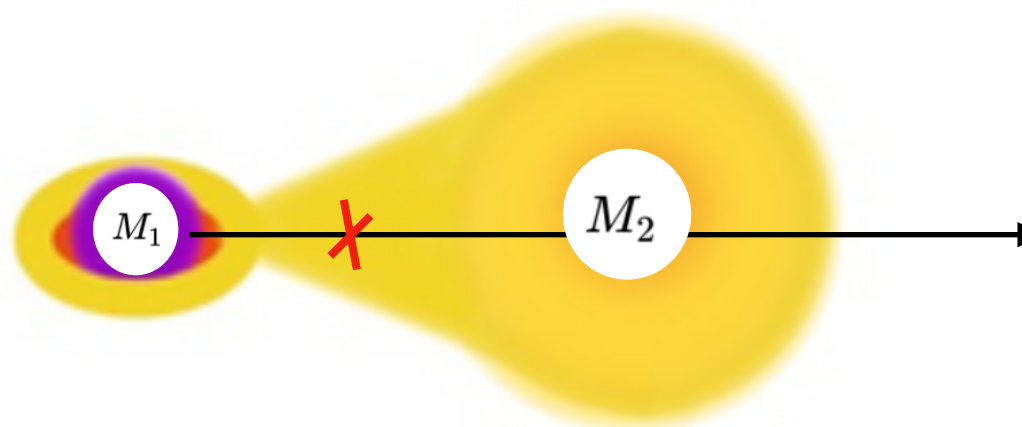


1. The function has asymptotes at $x = -a_1$ and $x = a_2$, and it tends to $-\infty$ at both sides of these asymptotes.
2. The function has three maxima which are called Lagrange points.
3. The function goes to $-\infty$ for $x \rightarrow \pm\infty$

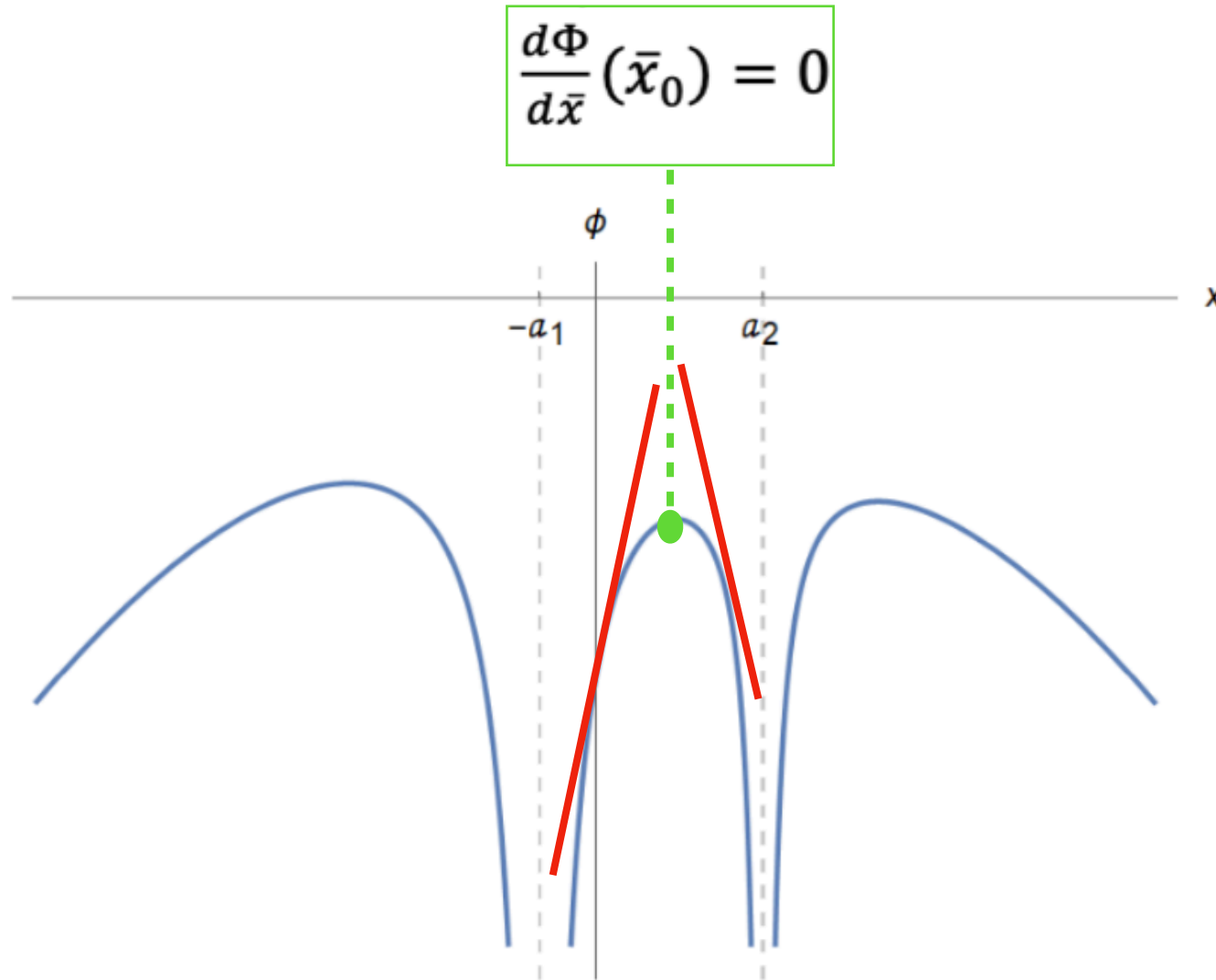
Question

Suppose $M_2 = M_1/3$ and assume that M_2 is surrounded by a rarefied gas of very low density. The mass of this gas is insignificant and we ignore its gravitational effects. If the size of this gas envelope becomes greater than a specific limit, the gas will overflow onto M_1 . Suppose the overflow occurs through $x = x_0$ on the x -axis.

A-3	Find the numerical value of $\frac{x_0}{a}$, up to two significant figures. You may use the calculator.	0.5 pt
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Solution



$$\Phi(\bar{x}, 0) = \frac{GM}{a} \left[-\frac{\frac{3}{4}}{\left(\bar{x} + \frac{1}{4}\right)} + \frac{\frac{1}{4}}{\left(\bar{x} - \frac{3}{4}\right)} - \frac{1}{2} \bar{x}^2 \right] \quad \bar{x} = x/a$$

$$f(\bar{x}) = \frac{a}{GM} \frac{d\Phi}{d\bar{x}}$$

$$\begin{aligned} &\begin{cases} f(0) > 0 \\ f(0.5) < 0 \end{cases} \rightarrow f(0.25) > 0 \rightarrow 0.25 < \bar{x}_0 < 0.5 \rightarrow f(0.375) < 0 \rightarrow \dots \rightarrow 0.358 < \bar{x}_0 < 0.361 \\ &\rightarrow f(0.360) > 0 \rightarrow 0.360 < \bar{x}_0 < 0.361 \rightarrow \boxed{\frac{x_0}{a} = \bar{x}_0 \approx 0.36} \end{aligned}$$

Question

Take the rotational period of the stars around their center of mass to be P . Assume that mass flows from M_2 to M_1 at a very small rate of $dM_1/dt = \beta$. This rate is so small that in each period of rotation, the distance between the two stars can be assumed to be constant. However, after a long period of time, the distance between the two stars changes, while the motion remains circular.

A.4. Calculate the rate of change of a and P in terms of β , M_1 , M_2 , and a

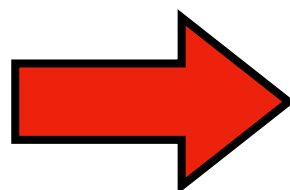
0.6 pt

Solution

$$J = \mu a V = \mu a^2 \omega = \frac{M_1 M_2}{M} a^2 \sqrt{\frac{GM}{a^3}} = \sqrt{\frac{GM_1^2 M_2^2}{M}} a$$

$$\frac{\dot{J}}{J} = \frac{\dot{M}_1}{M_1} + \frac{\dot{M}_2}{M_2} + \frac{1}{2} \frac{\dot{a}}{a}$$

$$\dot{M}_1 + \dot{M}_2 = 0$$



$$\frac{\dot{a}}{a} = -2 \frac{\dot{M}_1}{M_1} \left(1 - \frac{M_1}{M_2}\right) \rightarrow \dot{a} = -2\beta a \left(\frac{1}{M_1} - \frac{1}{M_2}\right)$$

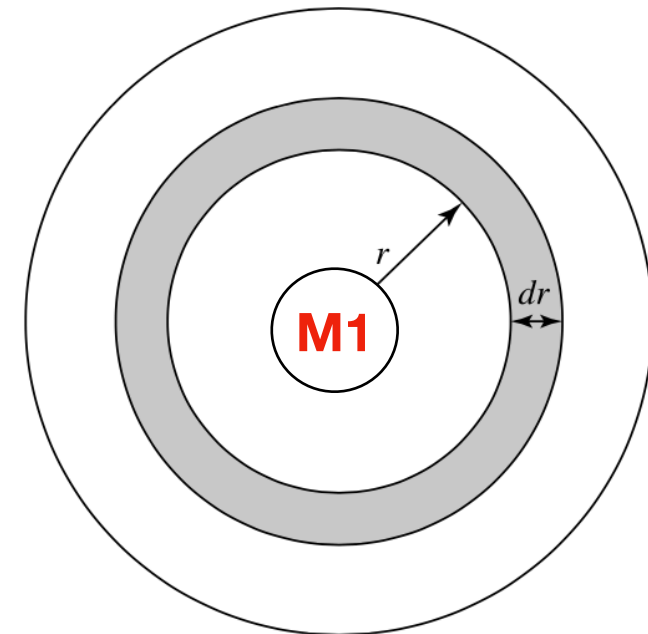
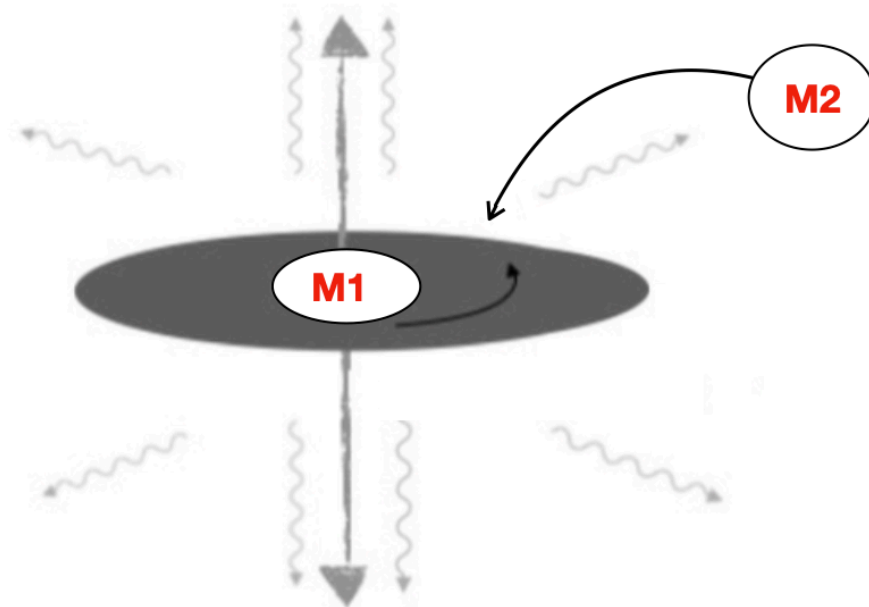
$$P = 2\pi \sqrt{\frac{a^3}{GM}} \rightarrow \frac{\dot{P}}{P} = \frac{3}{2} \frac{\dot{a}}{a} = -3 \frac{\dot{M}_1}{M_1} \left(1 - \frac{M_1}{M_2}\right) \rightarrow \dot{P} = -6\pi \sqrt{\frac{a^3}{GM}} \beta \left(\frac{1}{M_1} - \frac{1}{M_2}\right)$$

Question

The gas separated from M_2 forms a disk rotating around M_1 and heats up due to friction. As the gas loses energy, it spirals inward toward M_1 and finally falls onto it. In the steady state, the mass flows at the constant rate of β , from M_2 to the disk and from the disk onto M_1 . At the same time, the heated disk emits thermal radiation as a blackbody. This disk forms very close to the neutron star so the gravitational pull of the M_2 star can be ignored for the analysis of the disk's motion. Also ignore the heat capacity of the gas.

A.5. Determine the temperature of the disk at distance r from the center of the star M_1 in terms of β , M_1 , and physical constants.

1.0 pt



Solution

$$\frac{K(r) + U(r)}{m} = \frac{1}{2m}U(r) = \frac{-GM_1}{2r}$$

$$\frac{dE(r)}{dt dr} = \frac{-GM_1 \dot{m}}{2r} = \frac{GM_1 \beta}{2r}$$

$$\frac{dE(r)}{dt} = \sigma T(r)^4 2 \times 2\pi r dr = \frac{GM_1 \beta}{2r^2} dr$$

$$T(r) = \left(\frac{GM_1 \beta}{8\pi\sigma r^3} \right)^{1/4}$$

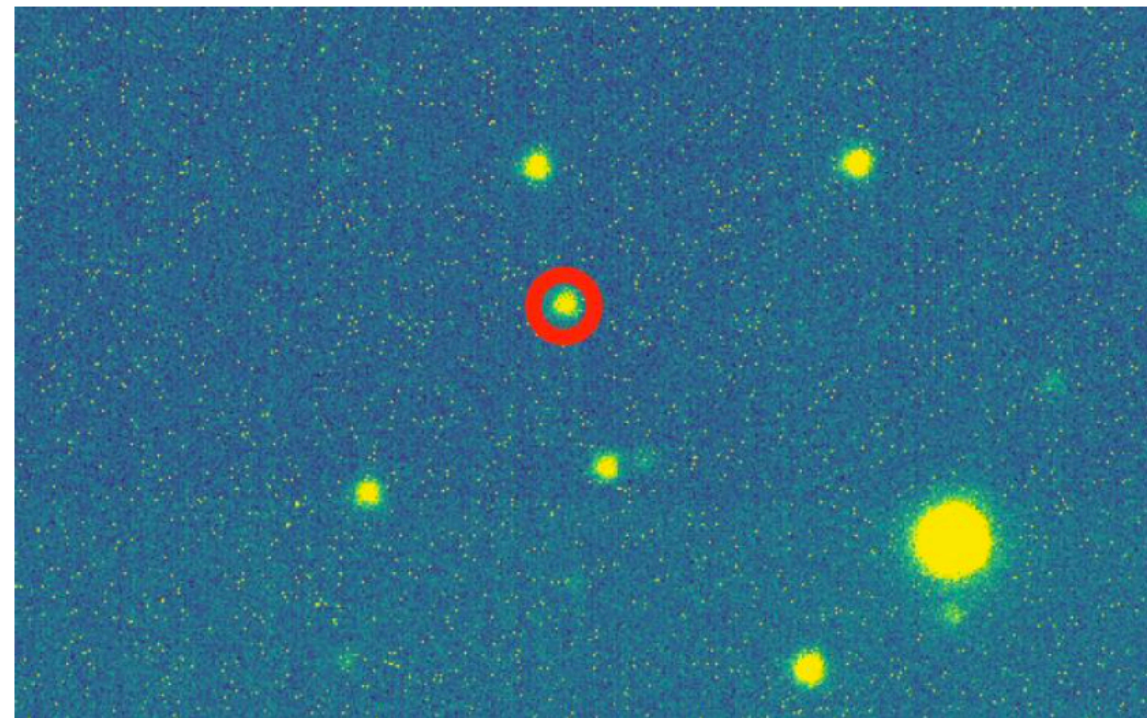
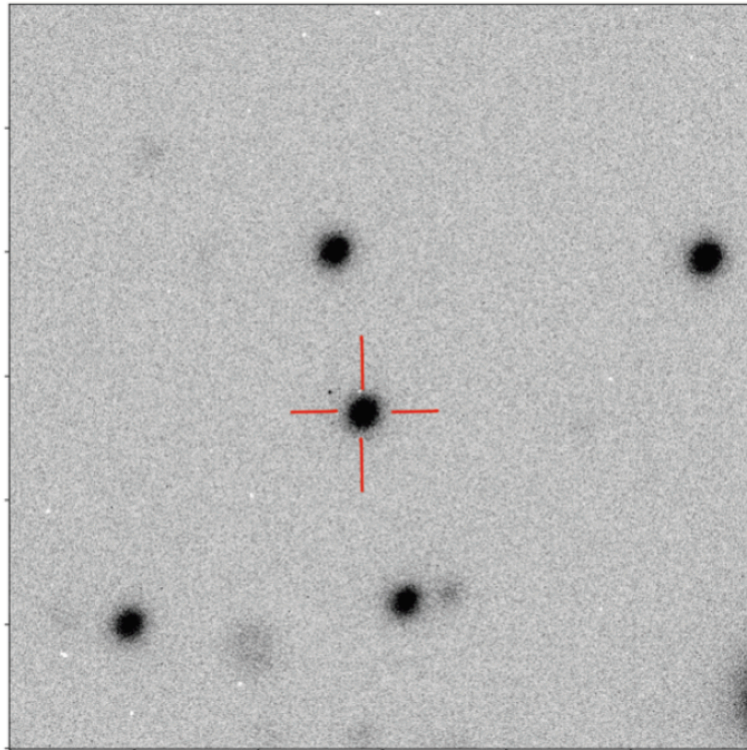


Figure 1 - The companion of the neutron star PSR J2215+5135

Question

In the binary system PSR J2215+5135, the mass of the neutron star is $M_{\text{NS}} = 2.27 M_{\odot}$ and the mass of its companion star is $M_S = 0.33 M_{\odot}$, where $M_{\odot} = 1.98 \times 10^{30} \text{ kg}$ is the mass of the Sun. The rotational period is $P = 4.14 \text{ h}$, and the Stefan-Boltzmann constant is $\sigma = 5.67 \times 10^{-8} \text{ W/m}^2\text{K}^4$, and the gravitational constant is $G = 6.67 \times 10^{-11} \text{ m}^3/\text{kg s}^2$. Assume that the mass flow rate to the neutron star is $\beta = \dot{M}_{\text{NS}} = 9 \times 10^{-10} M_{\odot} \text{ yr}^{-1}$.

A-6	Calculate the temperature of the disc at the radius $r = \frac{a}{10}$ in kelvins.	0.5 pt
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Solution

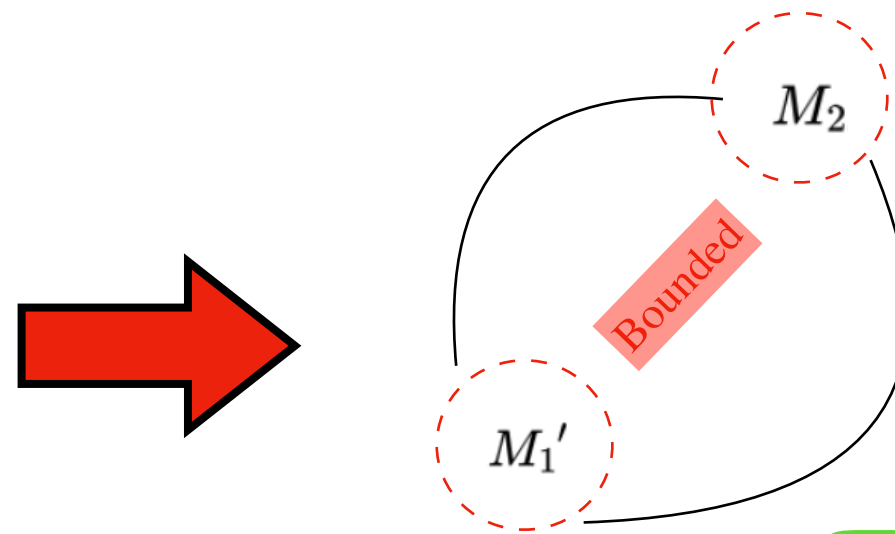
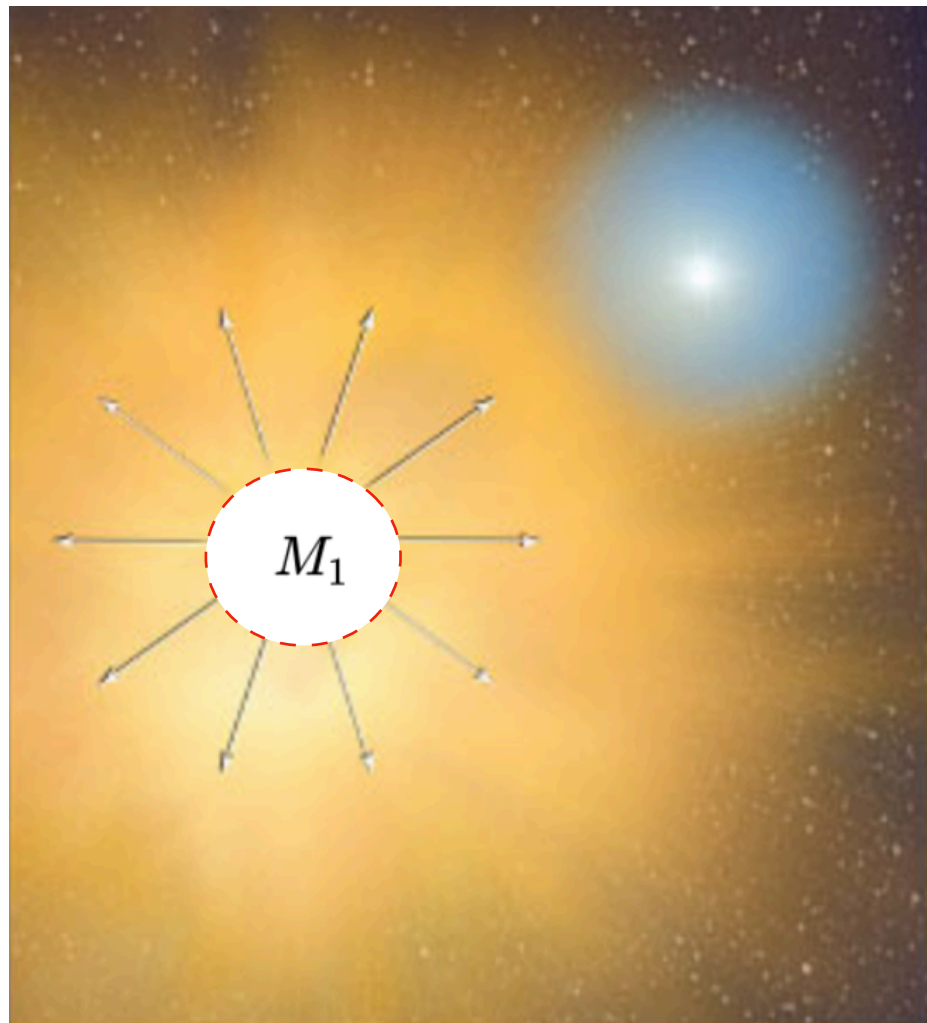
$$a = \left[\frac{P^2 G (M_S + M_{\text{NS}})}{4\pi^2} \right]^{\frac{1}{3}} = 1.2456 \times 10^9 \text{ m}$$

$$T = \left(\frac{G M_{\text{NS}} \beta}{8\pi \sigma r^3} \right)^{\frac{1}{4}} = 9 \times 10^3 \text{ K}$$

Question

Assume that after a sudden explosion, the M_1 star ejects a part of its mass out of the binary system at a very high speed, and its mass becomes M_1' . Take the magnitude of the velocity of M_1' relative to M_2 to be v' after the explosion.

A-7	Determine the maximum value of v' , in terms of M_1' , M_2 , and a , that allows the new binary system to stay bounded. Assuming that the explosion is isotropic, what is the minimum value of M_1' for the binary system to remain bounded?	0.7 pt
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Solution

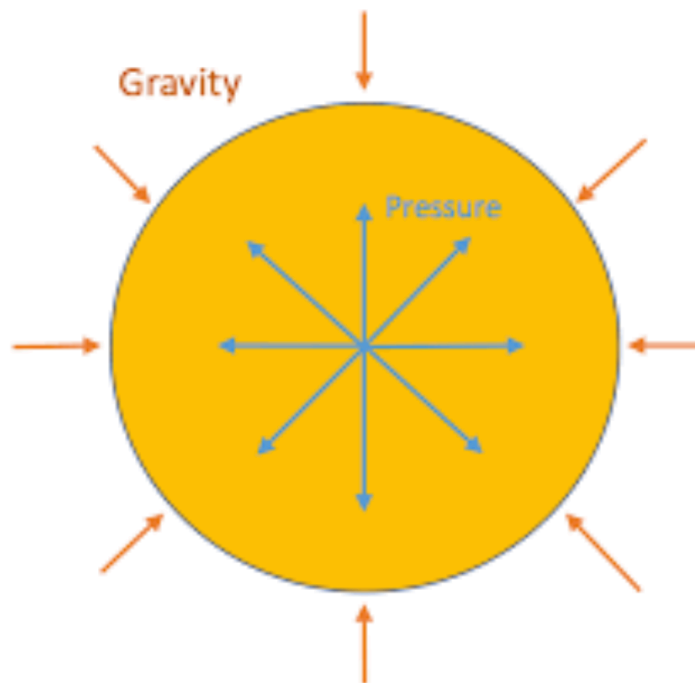
$$E' = \frac{1}{2}\mu'v'^2 - \frac{GM_1'M_2}{a} < 0 \rightarrow v' < \sqrt{\frac{2G(M_1' + M_2)}{a}}$$

isotropic explosion $v' = v = \sqrt{\frac{GM}{a}}$

$$\frac{M_1 - M_2}{2} < M_1'$$

$$\sqrt{\frac{G(M_1 + M_2)}{a}} < \sqrt{\frac{2G(M_1' + M_2)}{a}}$$

Analysis of the Stability of a Star



Question

In this part we study the stability of a single star. Consider a star containing a specific kind of matter with the equation of state $p = K\rho^\gamma$ where K and γ are constants. Let $p(r)$ and $\rho(r)$ be the pressure and density at a distance r from the center of the star, respectively. The pressure and density at the center of the star are p_c and ρ_c , respectively.

B.1. Determine the gravitational acceleration $g(r)$ near the center of the star in terms of r and the constants G and ρ_c .

0.2 pt

Solution

$$g = -\frac{4\pi G \int_0^r r'^2 \rho dr'}{r^2} \stackrel{\rho \cong \rho_c}{\cong} -\frac{4\pi G \rho_c r}{3}$$

Question

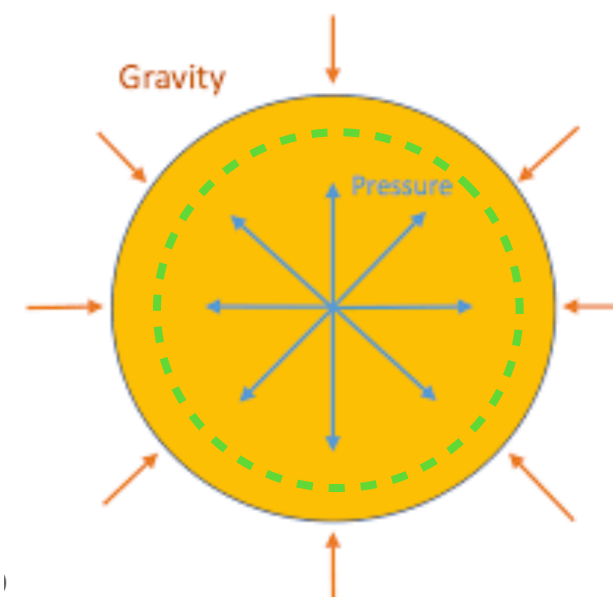
B.2. Derive a (differential) equation for determining $\rho(r)$ at equilibrium in the form

$\frac{d}{dr} \left[h_1(\rho, r) \frac{d\rho}{dr} \right] + h_2(r) \rho = 0$. Determine the functions h_2 and h_1 .

0.6 pt

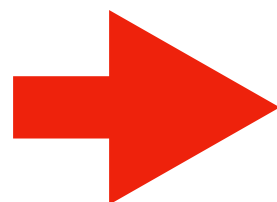
$$\vec{F} = -\frac{GM(\vec{r})\rho}{r^2} A \Delta r - \Delta p A = 0$$

Solution



$$p = K\rho^\gamma$$

$$M(r) = \int_0^r 4\pi r'^2 \rho(r') dr'$$



$$\frac{d}{dr} \left[r^2 \rho^{\gamma-2} \frac{d\rho}{dr} \right] + \frac{4\pi G r^2}{K\gamma} \rho(r) = 0$$

B.3. Construct a quantity r_0 of the form $r_0 = G^l p_c^m \rho_c^n$ with the dimension of length Question

0.4 pt

Solution

$$r_0 = G^{-\frac{1}{2}} p_c^{\frac{1}{2}} \rho_c^{-1}$$

B.4. Write the (differential) equation of task B.2 in the form $\frac{d}{dx} \left[j_1(u, x) \frac{du}{dx} \right] + j_2(x)u(x) = 0$ where $x = \frac{r}{r_0}$ and $u = \frac{\rho}{\rho_c}$. Question

0.3 pt

Solution

$$\frac{d}{dx} \left[x^2 u^{\gamma-2} \frac{du}{dx} \right] + \frac{4\pi x^2}{\gamma} u = 0$$

B.5. For $\gamma = 2$ one finds $u(x) = \frac{f(x)}{x}$. Determine $f(x)$.

Question

0.6 pt

Solution

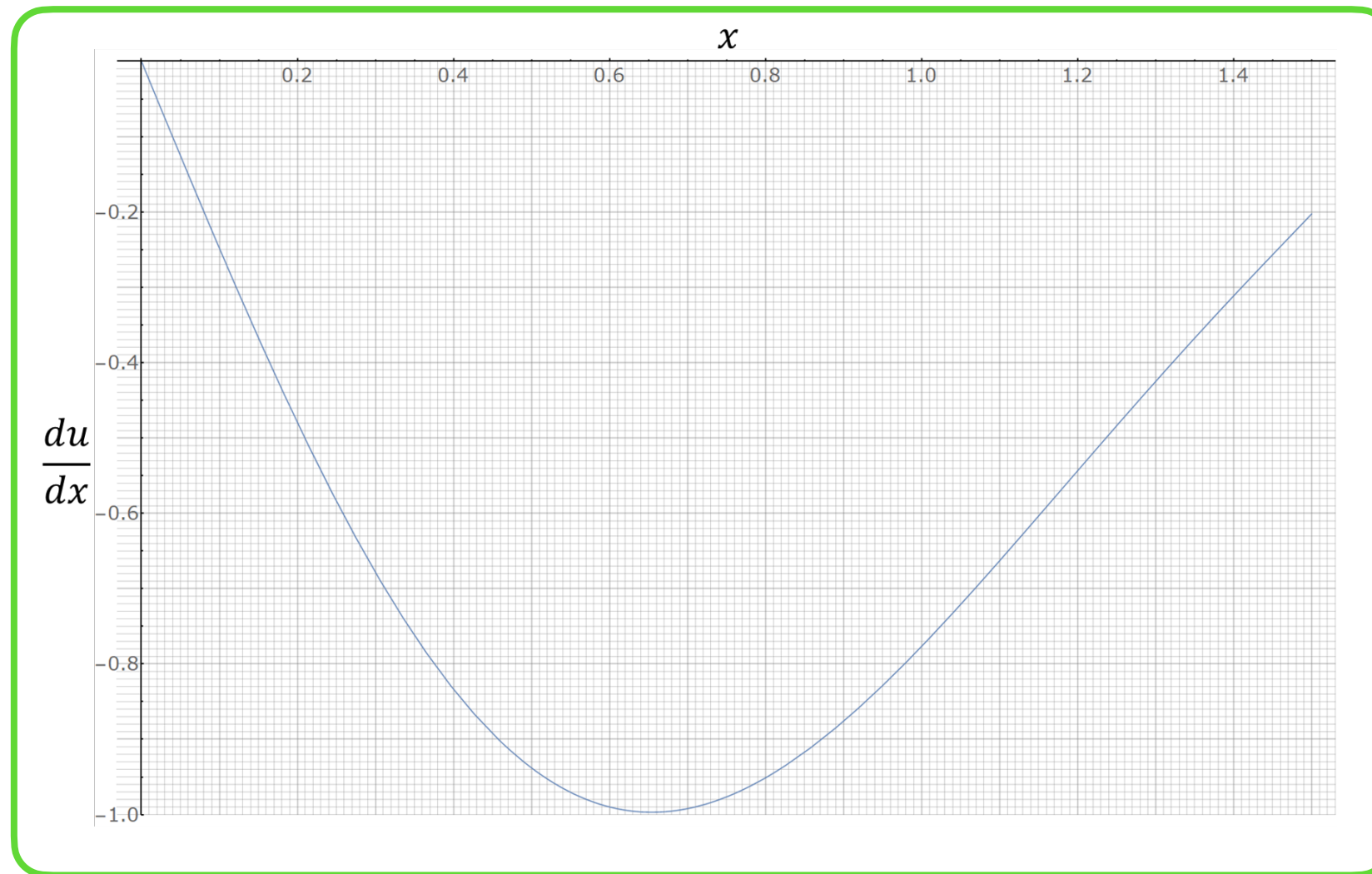
$$\gamma = 2 \rightarrow \frac{d}{dx} \left[x^2 \frac{du}{dx} \right] = -2\pi x^2 u(x) \rightarrow f''(x) = -2\pi f(x) \rightarrow f(x) = \frac{\sin(\sqrt{2\pi}x)}{\sqrt{2\pi}}$$

Question

Assume that for a particular star, $\frac{du}{dx}$ as a function of x , is represented by the curve given in Figure 2.

B.6. Using the behavior of the curve in Figure 2, in the vicinity of the point $x = 0$ find γ up to 3 significant figures. Use the given ruler if necessary.

0.8 pt

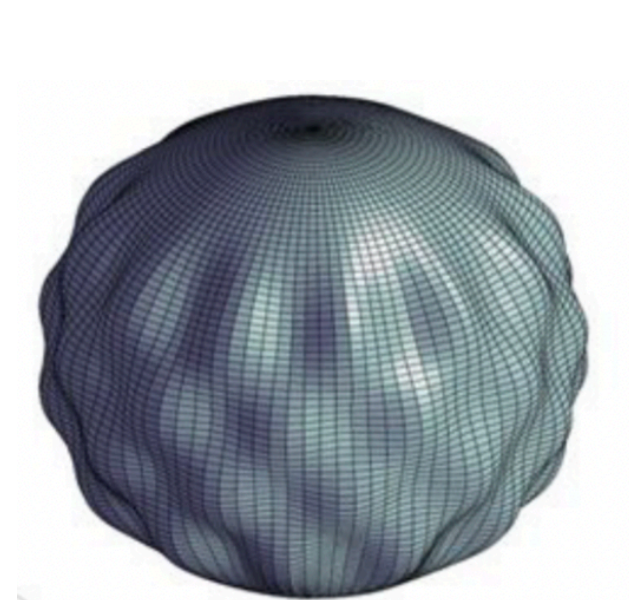
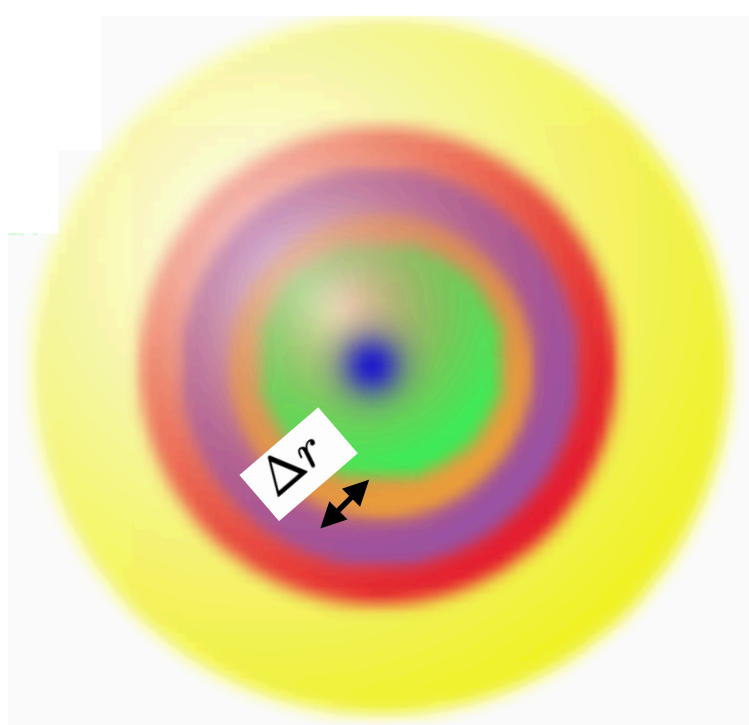


Solution

$$u''(0) + 2u''(0) + \frac{4\pi}{\gamma} = 0 \rightarrow \gamma = -\frac{4\pi}{3u''(0)}$$

$$\gamma \sim 1.66$$

$$u'(0) = 0, \quad \lim_{x \rightarrow 0} \frac{u'(x)}{x} = u''(0)$$

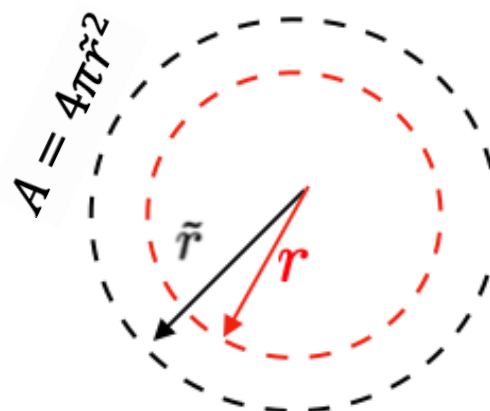


To analyze the stability of the system, we assume that the star deviates slightly from its equilibrium state: we assume that the spherical shell, which was in equilibrium at radius r , now has a radius \tilde{r} , similarly the parameters g , p , and ρ have changed to \tilde{g} , \tilde{p} , and $\tilde{\rho}$ respectively. For convenience, we shall only consider small r 's near the center of the star, for which we can assume that $\tilde{r} = r(1 + \varepsilon(t))$, where $\varepsilon(t) \ll 1$.

Question

B-7	Find $\tilde{\rho}$ and \tilde{g} in terms of ρ and g to the first order in ε .	0.9 pt
B-8	Using Newton's equation of motion for the spherical layer with the equilibrium radius of r find $\frac{d^2\tilde{r}}{dt^2}$ in terms of \tilde{g} , $\tilde{\rho}$, and $\frac{\partial\tilde{\rho}}{\partial\tilde{r}}$. (By $\frac{\partial\tilde{\rho}}{\partial\tilde{r}}$ we mean derivative of $\tilde{\rho}$ with respect to \tilde{r} at constant t .)	0.6 pt

$$\tilde{r} = r(1 + \epsilon(t))$$



Solution

B.7.

$$4\pi r^2 \rho(r) = 4\pi \tilde{r}^2 \tilde{\rho}(\tilde{r}, t) \frac{\partial \tilde{r}}{\partial r} \rightarrow \frac{\tilde{\rho}}{\rho} = \frac{r^2}{\tilde{r}^2} \left(\frac{\partial \tilde{r}}{\partial r} \right)^{-1} = (1 + \epsilon)^{-3} \cong 1 - 3\epsilon$$

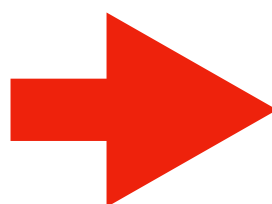
$$\frac{\tilde{g}}{g} = \frac{\frac{GM}{\tilde{r}^2}}{\frac{GM}{r^2}} = \frac{\frac{1}{\tilde{r}^2}}{\frac{1}{r^2}} = (1 + \epsilon)^{-2} \cong 1 - 2\epsilon$$

B-8

$$\vec{F} = -\frac{GM(\vec{r})\rho}{r^2} A \Delta r - \Delta p A$$

$$dM \ddot{\tilde{r}} = dM \tilde{g} - A d\tilde{p}$$

$$\frac{\partial \tilde{p}}{\partial \tilde{r}} = \tilde{\rho}(\tilde{g} - \ddot{\tilde{r}})$$



$$\ddot{\tilde{r}} = \tilde{g} - \frac{\left(\frac{\partial \tilde{p}}{\partial \tilde{r}} \right)}{\tilde{\rho}} = \tilde{g} - k\gamma \tilde{\rho}^{\gamma-2} \frac{\partial \tilde{\rho}}{\partial \tilde{r}}$$

Question

B-9	Obtain $\frac{d^2\epsilon}{dt^2}$ in terms of ϵ and the constants given in the problem. Find the minimum value of γ for a stable equilibrium, and find the oscillation's angular frequency of the star.	0.6 pt
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Solution

$$\tilde{r} = r(1 + \epsilon(t))$$

$$\tilde{\rho} \simeq \rho(1 - 3\epsilon)$$

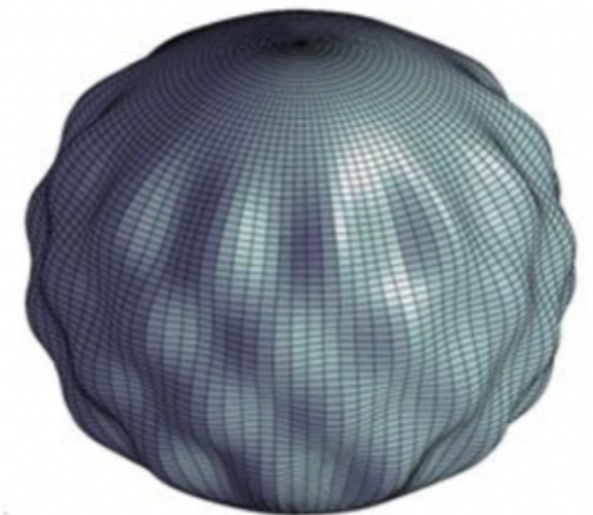
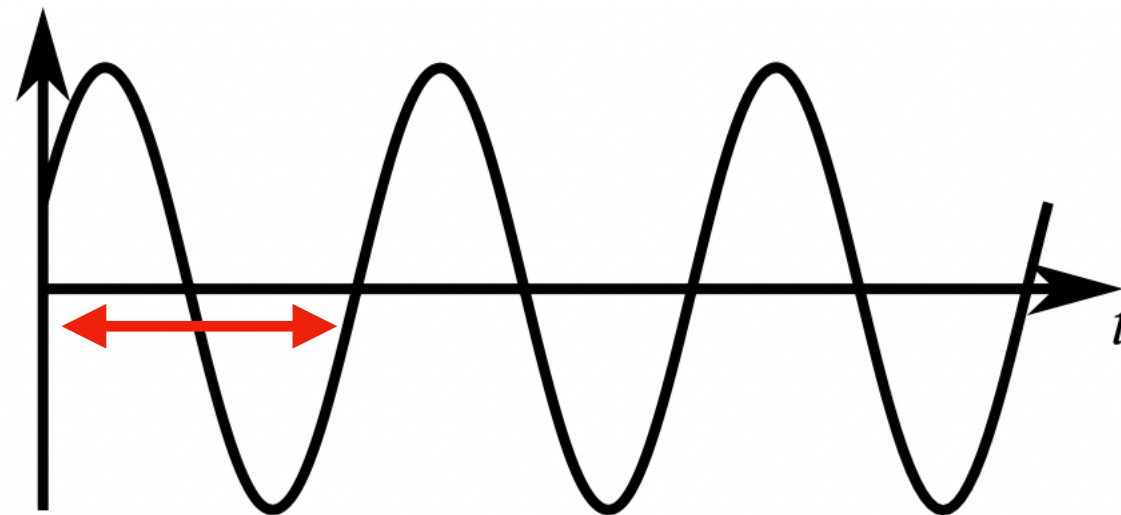
$$\tilde{g} \simeq g(1 - 2\epsilon)$$

$$\ddot{\epsilon} = -\frac{4\pi G \rho_c}{3}(3\gamma - 4)\epsilon$$

$$3\gamma - 4 > 0 \Rightarrow \gamma > \frac{4}{3}$$

$$\omega = \sqrt{\frac{4\pi G \rho_c}{3}(3\gamma - 4)}$$

ϵ



Syllabus

A1, A2, A3, A4, A6, A7

Mechanics, Kinematics, Statics, Dynamics, Celestial mechanics

A5

Thermodynamics and statistical physics (Black Body) + Mechanics

B1, B2, B3, B4 , B5, B6, B7

Mechanics, Dimension analysis, Mathematics, Calculus

B8,B9

Oscillations and Waves Harmonic oscillations and the equation of Harmonic Oscillation.

THANK YOU
FOR YOUR ATTENTION

ISFAHAN-IRAN