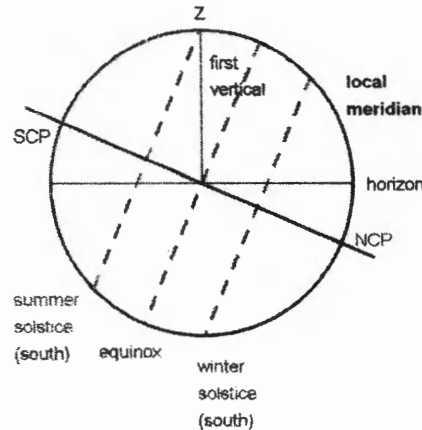


Theory - General Solutions

- 1 As the dial is vertical, the rod lies on the east-zenith-west plane. The clock works only when the Sun is in the northern hemisphere, $\delta_{Sun} > \delta_{zenith}$, because only in this case there will be a shadow projected over the dial.



- i) When $\delta_{Sol} \leq 0^\circ$, the Sun does not project a shadow when it is near the horizon (after rising and before setting). It corresponds to the time between the equinoxes: between September and March, during spring and summer.
- ii) When $\delta_{Sun} = < -22^\circ 54'$, the solar clock doesn't work at all. It corresponds to the small time interval close to the Summer solstice, during December and January.

3 pt
declination

2 pt
months & seasons

3 pt
declination

2 pt
months & seasons

- 2 Let be T_{sol} the duration of the solar day, T_{sid} the duration of the sidereal day and T_{year} the length of the year. In the present situation, we have

$$\frac{1}{T_{sol}} = \frac{1}{T_{sid}} - \frac{1}{T_{year}}$$

3 pt
calc. T_{sid}

Using the given data, $T_{sid} = 0.997270 \text{ days} = 23 \text{ h } 56 \text{ m } 4 \text{ s}$

3 pt
sidereal doesn't change

The sidereal day depends only of the magnitude of Earth's rotation speed, and the angular velocity of the sun's annual movement is opposite. So we have

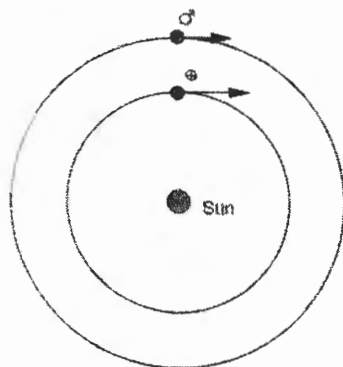
$$\frac{1}{T'_{sol}} = \frac{1}{T_{sid}} + \frac{1}{T_{year}}$$

2 pt
- to +

$$\therefore T'_{sol} = 0.994555 \text{ d} = 23 \text{ h } 52 \text{ m } 9.5 \text{ s}$$

2 pt
right number

3.



From Kepler's Third Law, we have:

$$\frac{a_J^3}{T_J^2} = \frac{a_M^3}{T_M^2} \Rightarrow T_J = 1,874 \text{ year}$$

5 pt
calculating the
period of Mars

Thus

$$\frac{1}{T_{op}} = \frac{1}{T_M} - \frac{1}{T_J} \Rightarrow T_{op} = 2,14 \text{ year}$$

5 pt
synodic period

4. The Moon's flux is proportional to its albedo, so $\frac{F_1}{a} = \frac{F_2}{1}$, being F_1 its actual flux, a the albedo and F_2 the hypothetical flux in the case the Moon had an albedo 1.
Being m_1 the Moon apparent magnitude and m_2 the desired value, we have:

$$m_2 - m_1 = -2,5 \log \frac{F_2}{F_1} \Rightarrow m_2 = m_1 + 2,5 \log a$$

5 pt
flux relation

Replacing $m_1 = -12,74$ and $a = 0,14$ we have: $m_2 = -14,9$

5 pt
all the rest

5. Fact 1: $\Delta\phi = 1^\circ = \frac{\pi}{180}$
 $L = 111 \text{ km}$

$$R_\oplus = \frac{L}{\Delta\phi}$$

2 pt
relations of fact 1

Fact 2: $g = 9,8 \text{ m/s}^2$

$$g = \frac{GM_\oplus}{R_\oplus^2} = \frac{4\pi G\rho_\oplus R_\oplus}{3} \therefore \rho_\oplus = \frac{3g\Delta\phi}{4\pi GL}$$

2 pt
relations of fact 2

Fact 3: $T_\oplus = 1 \text{ year}$

$$\frac{T_\oplus^2}{a_\oplus^3} = \frac{4\pi^2}{GM_\oplus} \therefore T_\oplus^2 = \frac{3\pi a_\oplus^3}{G\rho_\oplus R_\oplus^3}$$

$$M_\oplus = \frac{4}{3}\pi\rho_\oplus R_\oplus^3$$

2 pt
relations of fact 3

Fact 4: $\alpha = 30'$

$$\frac{\alpha}{2} \approx \frac{R_\oplus}{a_\oplus} \therefore T_\oplus^2 = \frac{3\pi a_\oplus^3}{G\rho_\oplus R_\oplus^3} = \frac{24\pi}{G\rho_\oplus \alpha^3}$$

2 pt
relations of fact 4

$$\therefore \frac{\rho_\oplus}{\rho_\odot} = \frac{\frac{3g\Delta\phi}{4\pi GL}}{\frac{24\pi}{GT_\oplus^2 \alpha^3}} = \frac{g\Delta\phi \alpha^3 T_\oplus^2}{32\pi^2 L} = 3,23$$

2 pt
answer

6. The reaction chain is



So the energy is

$$E = (2m_{\text{He}^3} - m_{\text{He}^4} - 2m_{\text{H}^1}) \cdot c^2 = 12,66 \text{ MeV}$$

7 pt
energy

$$\frac{\Delta m}{m} = \frac{12,66 \text{ MeV}/c^2}{2m_{\text{He}^3}} = 0,002254$$

3 pt
percentage

7. $T_1 = 30000 \text{ K}$

4 pt

$$T_z = 5000 \text{ K}$$

$$m_{bol} = \text{constant} \Leftrightarrow L = \text{constant}$$

As $L = 4\pi\sigma R^2 T^4$ and $L_1 = L_2$, we have:

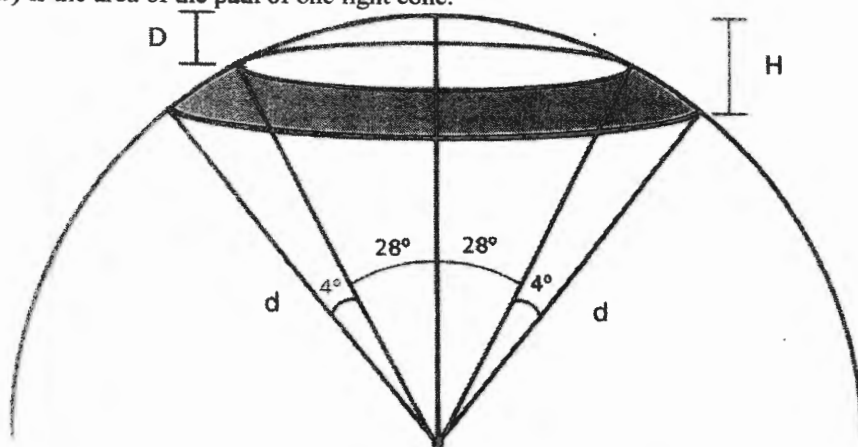
$$4\pi\sigma R_0^2 T_1^4 = 4\pi\sigma R_{shell}^2 T_2^4 \Rightarrow \frac{R_{shell}}{R_0} = 36$$

8. As for the probability, supposing that the neutron star space orientation is random, is

$$p = \frac{2A_p(d)}{4\pi d^2}$$

2 pt
stating p

where $A_p(d)$ is the area of the path of one light cone.



3 pt
evaluating p

As the figure above shows, the area is the difference of the spherical caps of height H and D:

$$A = 2\pi d(H - D) = 2\pi d^2 \left[\cos\left(\theta - \frac{\alpha}{2}\right) - \cos\left(\theta + \frac{\alpha}{2}\right) \right]$$

$$p = \cos\left(\theta - \frac{\alpha}{2}\right) - \cos\left(\theta + \frac{\alpha}{2}\right) = 3.5 \cdot 10^{-2} \text{ or } 3.5\%$$

2 pt
stating the flux

The light waves are spread along the double cone and not along a sphere, thus:

$$F = \frac{L}{2A(d)} = \frac{10\,000 L_\odot}{4\pi d^2 \left(1 - \cos\frac{\alpha}{2}\right)}$$

Where $A(d)$ is the area of one light cone. Using the Sun's absolute magnitude, we can find the absolute magnitude of the pulsar (m_p):

3 pt
evaluating
magnitude

$$m_p - M_\odot = -2.5 \log\left(\frac{F}{\frac{L_\odot}{4\pi(10\text{ pc})^2}}\right) = -2.5 \log\left(10\,000 \cdot \frac{1}{1 - \cos\frac{\alpha}{2}} \left(\frac{10}{1000}\right)^2\right)$$

$$m_p = -3.24$$

9. The intrinsic color indexes are $U - V = 0.300$ and $U - V = 0.330$. therefore
 $U - B = (U - V) - (B - V) = 0.030$

The difference between the real and the measured color indexes are due to the absorption in the interstellar medium, as well as due the stronger absorption of the planetary nebula (in fact, during the calculations, one can easily see that the interstellar absorption is not enough to account for the measured values). So, for the first (nearest) star:

$$\begin{aligned}(B - V)_1 &= (B - V) + (E_B - E_V) \cdot D + (E'_B - E'_V) \cdot R = 0.327 \\(U - V)_1 &= (U - V) + (E_U - E_V) \cdot D + (E'_U - E'_V) \cdot R = 0.038 \\(U - B)_1 &= (U - B) + (E_U - E_B) \cdot D + (E'_U - E'_B) \cdot R = 0.365\end{aligned}$$

3 pt
equation

apparentintrinsic interstellar nebulaabsorption
absorption

Calculating the interstellar absorption, we can find for, the nebula values:

$$\begin{aligned}(E'_B - E'_V) \cdot R &= 0.0155 \\(E'_U - E'_V) \cdot R &= 0.0100 \\(E'_U - E'_B) \cdot R &= -0.0055\end{aligned}$$

1 pt
absorption
values

So, for the second (farthest) star,

$$\begin{aligned}(B - V)_2 &= (B - V) + (E_B - E_V) \cdot 3D + (E'_B - E'_V) \cdot 2R = 0.3655 \\(U - V)_2 &= (U - V) + (E_U - E_V) \cdot 3D + (E'_U - E'_V) \cdot 2R = 0.425 \\(U - B)_2 &= (U - B) + (E_U - E_B) \cdot 3D + (E'_U - E'_B) \cdot 2R = 0.0595\end{aligned}$$

4 pt
equation

2 pt
final answers

10. As we have a differential equation that should be integrated, we expect that the student can solve it in two ways: using an approximation (Way 1) or calculus (Way 2).

Common

Way 1

Using the definition of Hubble constant: $\frac{da}{dt} = Ha$, where a is the scale factor and using the

approximation $\frac{da}{dt} \approx \frac{\Delta a}{\Delta t}$, and supposing that the variation of the scale factor and the Hubble constant is small compared to their values, and using that $aT = \text{const}$:

$$\frac{\Delta a}{\Delta t} = H_i a_i \Rightarrow \frac{\Delta T}{\Delta t} = -H_i T_i \Rightarrow \Delta t = -\frac{1}{H_i} \frac{\Delta T}{T_i} = 5.0 \cdot 10^8 \text{ years}$$

2 pt
Finding final
scale factor

Way 1
7 pt
Using the
approximation
correctly

1 pt
final answer

Way 2

As we have a matter dominated Universe ($\rho \propto a^{-3}$), without dark energy, and a density parameter $\Omega_0 = 1$, we have:

$$\Omega_0 = \frac{8\pi G \rho}{3H^2} = \frac{8\pi G \zeta}{3H^2 a^3} = 1$$

From the Hubble law, $\frac{da}{dt} = Ha$:

$$\left(\frac{da}{dt}\right)^2 = \frac{8\pi G \zeta}{3 a} \Rightarrow \frac{da}{dt} = \sqrt{\frac{8\pi G \zeta}{3 a}}$$

4 pt
finding the diff.
equation

Integrating this equation:

$$\frac{2}{3} \left[a_f^{3/2} - a_i^{3/2} \right] = \sqrt{\frac{8\pi G \zeta}{3}} \Delta t \Rightarrow \Delta t = \frac{2}{3} \left[\left(\frac{a_f}{a_i} \right)^{3/2} - 1 \right] \frac{1}{\sqrt{\frac{8\pi G \zeta}{3 a_i^3}}} = \frac{2}{3} \frac{1}{H_i} \left[\left(\frac{a_f}{a_i} \right)^{3/2} - 1 \right]$$

3 pt
Integrating
equation

Using the fact that $aT = \text{const}$:

$$a_i T_i = a_f T_f \Rightarrow \frac{a_f}{a_i} = \frac{T_i}{T_f} = \frac{2.73}{2.63} = 1.038$$

Obs:
(7pt for using
correctly that
 $a \propto t^{2/3}$)

So the time is:

$$\Delta t = 5.2 \cdot 10^8 \text{ years}$$

1 pt
final answer

- 11 The oscillation occurs due to the fact that the center of mass (CM) of the Solar System does not lie exactly at the center of the Sun, but in another position around which the center of the Sun revolves.

Given $a_{21} = 5.204 \text{ AU}$, the distance from CM to the center of the Sun is:

$$x_{CM} = \frac{\frac{1}{1047}}{1 + \frac{1}{1047}} 5.204 = 0.004966 \text{ AU},$$

3 pt
position of CM

Which is the amplitude of the movement.

If the distance to Barnard's Star is $d_{bar} = 1.834 \text{ pc}$, we can write, for the angular amplitude:

$$\alpha \approx \frac{x_{cm}}{d_{bar}} = 0.002708''$$

3 pt
angular
amplitude

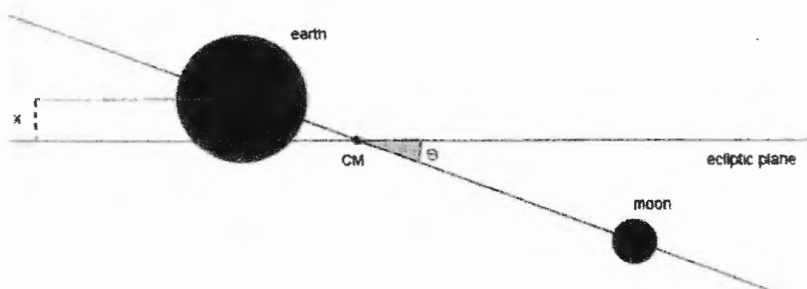
As for the period, it is the period of Jupiter's orbit:

$$\frac{T_{21}^2}{a_{21}^3} = \frac{T_{\oplus}^2}{a_{\oplus}^3} = 1 \frac{\text{year}^2}{\text{AU}^3} \therefore T_{Jup} = 11.87 \text{ years}$$

4 pt
period of
oscillation

Obs: the difference of Jupiter's mass in the Kepler's Third Law represents a difference in less than 0.01 year at the period of Jupiter.

- 12 The phenomenon is as follow:



The Center of Mass of the Earth-Moon system (CM) lies on the plane of the ecliptic, so as the Lagrangian points of its orbit around the Sun. Let be D the distance between Earth and Moon and x the distance between Earth and the CM. We then have:

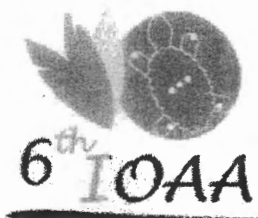
$$M_{\oplus} \cdot x = M_{\ell} \cdot (D - x) \therefore x = \left(\frac{M_{\ell}}{M_{\oplus} + M_{\ell}} \right) D$$

2 pt
equating the CM

The Earth's displacement perpendicular to the ecliptic is $A = 2x \sin \theta$, where $\theta = 5^\circ$ is the inclination of the plane of Moon's orbit relative to the plane of ecliptic. Calculating:

$$A = \frac{2D M_{\ell}}{M_{\oplus} + M_{\ell}} \sin 5^\circ = 810 \text{ km}$$

4 pt
calculating the
displacement



6th International Olympiad on Astronomy and Astrophysics

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3 pt
telescope
resolution

1 pt
answer

The telescope resolution is given by

$$r'' = 1.22 \frac{\lambda}{\delta} = \frac{A}{d}$$

where δ is the diameter of the objective lens and d is the distance to the telescope to Earth. But, at L4(L5), Earth and Sun forms an equilateral triangle, so $d = 1$ UA. Using $\lambda = 400$ nm (that is, the visible wavelength that allows the smallest telescope), we find $\delta = 9.01$ cm

- 13 As seen from an inertial reference frame outside Earth, the observer must describe a circle parallel to the ecliptic, with constant speed, in a period of one sidereal day. Since the ecliptic poles are in zero altitude, the circle will be a great circle.

But the resulting movement should be the actual movement of the observer (in relation of Earth's surface) plus the movement of Earth (in relation to the inertial reference frame). After some time t , the observer must describe an arc of length $\phi = 2\pi/T$ on the great circle; at the same time, the Earth will rotate by this same amount.

In other words, the displacement of the observer is the composition of a rotation with an angle ϕ around the ecliptic axis and a rotation with an angle $-\phi$ around the earth's rotation axis. But this is exactly the same construction which describes an analemma of a planet with a circular orbit, inclination of $\epsilon = 23.44^\circ$ between the ecliptic and equator!

So, the important characteristics of the drawing curve are:

- It is closed
- Maximum and minimum latitudes are $\theta = \pm 23.44^\circ$
- It is symmetric around the Equator
- It is symmetric around a meridian
- It has the shape of an eight (8)

1 pt
each item

For the calculation: as the observer starts on the southern hemisphere, he will cross the Equator for the first time in the direction south-north. Its velocity relative to an inertial referential will have intensity $v = 2\pi R/T$ and azimuth (the angle with the north-south direction) $\alpha = 90^\circ - \epsilon$. The velocity of the Earth's surface will have the same intensity and azimuth $\alpha = 90^\circ$. Subtracting the vectors, we'll have a resultant vector with modulus

$$v' = v\sqrt{2(1 - \cos \epsilon)}$$

And azimuth

$$\alpha' = -\frac{\epsilon}{2} = -11.72^\circ$$

3 pt
velocity

2 pt
direction

14. Let be R the radius of the Celestial Sphere and (x, y, z) a coordinate system with origin at the observer, z is the down-up axis, y is the east-west axis and x is the south-north axis.

In the superior culmination, the altitude angle is $h = 7.41 + 49.35 = 56.76^\circ$ and we have

$$\vec{v} = v_0 \hat{y} = 2\pi R \hat{y} / \text{day}$$

$$\vec{r}_0 = R(-\cos h \hat{x} + \sin h \hat{z})$$

$$\mathbf{r}(t) = \mathbf{r}_0 + \mathbf{v}t$$

The altitude in function of our coordinate system is $h = \tan^{-1} \left(\frac{z}{\sqrt{x^2 + y^2}} \right)$. When $t \rightarrow \infty$, the coordinates x and z of the star remain constant, while $y \rightarrow \infty$, and we have

$$h = \tan^{-1} \left(\frac{z}{\sqrt{x^2 + y^2}} \right) \xrightarrow{t \rightarrow \infty} \tan^{-1} 0 = 0^\circ$$

Similarly, because $y \rightarrow \infty$ but x and z remain constant, the final azimuth is the West direction ($A = 90^\circ$ or $A = 270^\circ$, depending on the system used).

For the magnitude, $m_0 - m_6 = -2.5 \log \frac{F_0}{F_6}$ and $F = \frac{L}{4\pi R^2}$ so $R_6 = 12.88 R_0$

$$|\mathbf{r}(t)| = \sqrt{x^2 + y^2 + z^2} = R\sqrt{1 + 4\pi^2 t^2} = 12.88 R$$

$$t = 2.044 \text{ days}$$

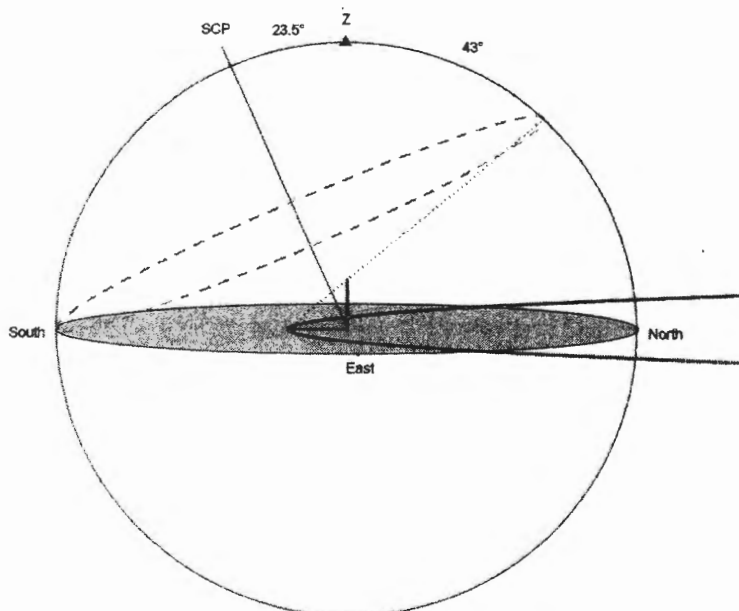
15. A known geometrical result is that the projection of a circle in an arbitrary plane is a conic curve (afterwards, projecting a circle is drawing a cone that intersects the given plane). In this particular case, the Sun crosses the horizon at one point, when the tip of the shadow goes to infinity. So the projection must be a parabola.

The equation of the curve will have the form $y = ax^2 + bx + c$. We just need to find the coefficients of this curve.

To find c , we can take the Sun in its maximum culmination, lying in the North-South plane. In this case, the altitude angle will be $23^\circ 27' + 23^\circ 27' = 46^\circ 54'$ and, as $x = 0$.

$$y = c = -39.6 \cdot \cotan 46.9 = -37.06$$

To find b , we just need to notice that the parabola must be symmetrical around the y axis, so $b = 0$.



3 pt
 $h \rightarrow 0$

2 pt
 $A \rightarrow W$

2 pt
ratio of distances

3 pt
time

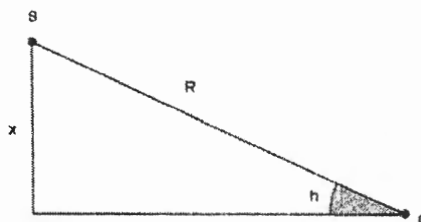
3 pt
it's a parabola

1 pt
equation of parabola

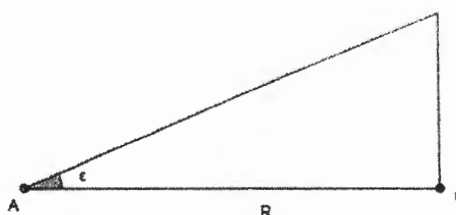
2 pt
finding c

1 pt
finding b

To find a , we can look at the Sun when it is crossing the First Vertical or the East-West plane. Let's call P this point. Looking at the E-W plane, we have the triangle on the side, where O is the origin of the system of coordinates, h is the altitude angle, R is the radius of the celestial sphere and x is the height of the Sun. We know that the altitude angle is such as $\sin h = x/R$.



On the other hand, looking at the North-South plane, we have another triangle. O is the origin, R is the radius of the celestial sphere (lying in the plane of horizon), A is the point of inferior culmination of the Sun (when it touches the horizon), ϵ is the obliquity of ecliptic, the hypotenuse lies on the plane containing the Sun trajectory during that day, and x lies in the E-W plane being the height of the Sun as defined above.



3 pt
finding a

Then.

$$x = R \tan \epsilon \quad \therefore \quad \sin h = \tan \epsilon \quad \therefore \quad h = 25^\circ 42'$$

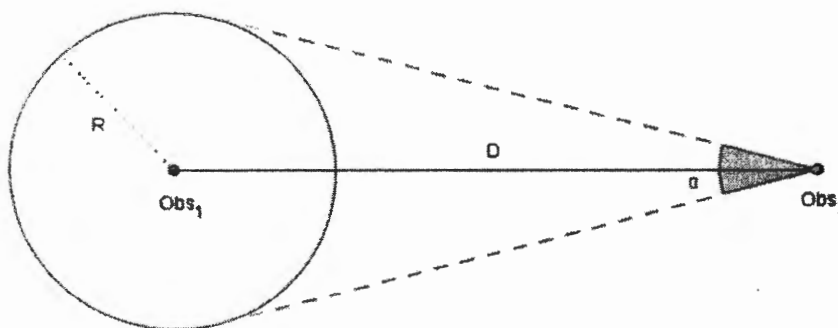
So, when $y = 0$, $x = 39.6 \cdot \cotan 25.7 = 82.28$ and

$$a = -\frac{c}{x^2} = 0.0055$$

And the equation is:

$$y = 0.0055x^2 - 37.1$$

A. Look at the figure:



$$\tan \alpha \approx \alpha \approx \sin \alpha \approx \frac{2R}{D}$$

$$\Delta m = m_B - m_A = -2.5 \log \frac{F_B}{F_A}$$

The brightness as seen by the astronomer (F_A) is the brightness of N stars with absolute magnitude

M_0 (and luminosity L_0) at a distance D from him:

$$F_A = N \frac{L_0}{4\pi D^2}$$

We assume the cluster is homogeneous, so the density of stars ρ inside the cluster should be constant. The brightness coming from a spherical shell with small thickness ΔR at a distance R' is

$$F_{shell}(R', \Delta R) = \rho V_{shell} \frac{L_0}{4\pi R'^2} = \frac{N}{\frac{4}{3}\pi R^3} 4\pi R'^2 \Delta R \frac{L_0}{4\pi R'^2} = \frac{3NL_0\Delta R}{4\pi R^3}$$

The brightness doesn't depend on the radius of the shell R' . If we divide the cluster in n shells, $R = n\Delta R$, the brightness of the cluster as seen by the astronomer 2 is:

$$F_B = nF_{shell} = n \frac{3NL_0\Delta R}{4\pi R^3} = \frac{3NL_0}{4\pi R^2}$$

So the difference in the magnitudes Δm is:

$$\Delta m = -2.5 \log \left(\frac{\frac{3NL_0}{4\pi R^2}}{\frac{3NL_0}{4\pi D^2}} \right) = -2.5 \log \left[3 \left(\frac{D}{R} \right)^2 \right] = 2.5 \log \left(\frac{\alpha^2}{12} \right)$$

A2.

The total energy received is proportional to the objective area:

$$F'_A = F_A \frac{\pi \left(\frac{x}{2} \right)^2}{\pi \left(\frac{d_{pup}}{2} \right)^2} \Rightarrow x = 6 \cdot 10^{-3} \sqrt{\frac{F'_A}{F_A}}$$

$$F'_A = F_B \Rightarrow \frac{F'_A}{F_A} = \frac{F_B}{F_A} = \frac{12}{\alpha^2} \Rightarrow x = 6 \cdot 10^{-3} \sqrt{\frac{12}{\alpha^2}} = \frac{0.021}{\alpha} \text{ m}$$

A3.

In this case, a fraction $\xi(d) = \frac{A(\alpha)}{4\pi d^2}$ of the light that comes from the distance d is seen by the biologist. $A(\alpha)$ is the area of the interior of a cone with aperture α that intersects the sphere with radius d and center in the vertices of the cone (a spherical cap of radius d and height D):

$$\cos \left(\frac{\alpha}{2} \right) = \frac{d-D}{d} \text{ and } A(\alpha) = 2\pi Dd \Rightarrow A(\alpha) = 2\pi d^2 \left[1 - \cos \left(\frac{\alpha}{2} \right) \right] \Rightarrow \xi = \sin^2 \left(\frac{\alpha}{4} \right)$$

So the fraction is independent of the distance d , and the fraction of light that the biologist sees is $F'_B = \xi F_B$:

$$\Delta m' = m'_B - m_A = -2.5 \log \frac{F'_B}{F_A} = -2.5 \log \left[\frac{12 \sin^2 \left(\frac{\alpha}{4} \right)}{\alpha^2} \right]$$

B. B1.

1st line: $\{(0.0); (15.210)\} (10^3 \text{ ly}, \frac{\text{km}}{\text{s}})$

2nd line: $\{(15.210); (60.200)\} (10^3 \text{ ly}, \frac{\text{km}}{\text{s}})$

4pt
stating the flux A

5pt
no-R'

4pt
stating the flux B

2 pt
calculate answer

3 pt
x as function of
Fs

3 pt
calculating the
diameter

2 pt
formulate the
fraction of flux

4 pt
calculate

3 pt
calculating the
new magnitude

4 pt
For the lines (2
pt each,
considered right
with some

$$V(D) = \begin{cases} 14D, & 0 \leq D < 15 \\ 213 - 0.22D, & 15 \leq D \leq 60 \end{cases}, D \text{ in } 10^3 \text{ l.y. and } V \text{ in km/s}$$

B2.

Using the equations from (a):

$$\omega(D) = \frac{V(D)}{D} \Rightarrow \omega(D) = \begin{cases} 14, & 0 \leq D < 15 \\ \frac{213}{D} - 0.22, & 15 \leq D \leq 60 \end{cases} \frac{\text{km/s}}{10^3 \text{ l.y.}}$$

$$\omega_{\text{wave}}(D) = 2\omega(D)$$

1 pt
writing omega

The time that a spiral arm takes to have one more turn around the galactic center:

$$\omega_{\text{rel}} = \omega_{\text{wave}}^{\text{max}} - \omega_{\text{wave}}^{\text{min}}$$

$$\omega_{\text{rel}} = \frac{1}{2} \left[\left(\frac{213}{15} - 0.22 \right) - \left(\frac{213}{80} - 0.22 \right) \right] = 23.9 \frac{\text{km/s}}{10^3 \text{ l.y.}}$$

$$P_{\text{sin}} = \frac{2\pi}{\omega_{\text{rel}}} = 8.15 \cdot 10^7 \text{ years}$$

4pt
relative velocity

1pt
answer

B3.

Tully-Fisher relation estimates the luminosity of the galaxy, and considering $\Delta V = 2V_{\text{max}}$:

$$L = 16kV_{\text{max}}^4$$

2 pt
Use Tully-Fisher
with Vmax

We can compare the magnitude of the galaxy with the Sun:

$$m - M = 5 \log d - 5 \Rightarrow m - \left(M_{\text{Sun}} - 2.5 \log \frac{L}{L_{\text{Sun}}} \right) = 5 \log d - 5$$

$$\log d = 1 + \frac{m - M_{\text{Sun}} + 2.5 \log \left(\frac{16kV_{\text{max}}^4}{L_{\text{Sun}}} \right)}{5}$$

2 pt
Compare with
the Sun

From the graph: $V_{\text{max}} = 230 \text{ km/s}$:

$$\log d = 1 + \frac{8.5 - 4.8 + 2.5 \log(16 \cdot 0.317 \cdot 230^4)}{5} = 6.816 \Rightarrow d = 6.5 \text{ Mpc}$$

2 pt
Answer (use of a
reasonable
V_max)

B4.

We can estimate the recession velocity of the galaxy using Hubble's law:

$$V_0 = H \cdot d$$

The maximum and minimum wavelengths occur where the radial velocity is V_{max} , the radial velocity is:

$$V = V_0 \pm V_{\text{max}}$$

Using $\frac{\Delta \lambda}{\lambda} = \frac{v}{c}$, the wavelengths are:

$$\lambda_{\text{max}} = \lambda \left(1 + \frac{Hd + V_{\text{max}}}{c} \right)$$

$$\lambda_{\text{min}} = \lambda \left(1 + \frac{Hd - V_{\text{max}}}{c} \right)$$

3pt
Write the
velocities
correctly

2 pt
Doppler law

Calculating the values, using $\lambda = 656.28 \text{ nm}$:

$$\lambda_{\text{max}} = 657.79 \text{ nm}$$

$$\lambda_{\text{min}} = 656.78 \text{ nm}$$

B5. Supposing a mass with spherical symmetry distribution and a circular orbit, we have:

$$\frac{V^2}{D} = \frac{GM_{int}}{D^2} \Rightarrow M_{int} = \frac{V^2 D}{G}$$

From the graph $V(3.0 \cdot 10^4 \text{ l.y.}) = 225 \text{ km/s}$, So the mass is:

$$M_{int} = 2.15 \cdot 10^{41} \text{ kg} = 1.08 \cdot 10^{11} M_{sun}$$

B6.

As the mass is derived from the rotation curve, it is necessary to take into account that part of the mass is dark matter, and only part of the mass is in the stars ($\rho_{star} = 1/3$). The fraction of baryonic matter is:

$$\rho_{matter} = \frac{f_{baryon}}{f_{matter}} = \frac{4}{22 + 4} = \frac{4}{26}$$

The number of stars is then:

$$N = \rho_{matter} \cdot \rho_{star} \frac{M_{galaxy}}{M_{Sun}} = \frac{4}{26} \frac{1}{3} 1.08 \cdot 10^{11} = 5.5 \cdot 10^9 \text{ stars}$$

5 pt
Estimate the
internal mass

2pt
take in account
the dark matter
correctly

2 pt
Calculate the
number